# The relaxed energy for $S^{2}$-valued maps and measurable weights 

# L'énergie relaxée pour des applications à valeurs dans $S^{2}$ et poids mesurables 

Vincent Millot<br>Laboratoire J.L. Lions, Université Pierre et Marie Curie, B.C. 187<br>4 Place Jussieu, 75252 Paris Cedex 05, France<br>E-mail adress: millot@ann.jussieu.fr


#### Abstract

We compute explicitly a relaxed type energy for maps $u: \Omega \subset \mathbb{R}^{3} \rightarrow S^{2}$. The explicit formula involves the length of a minimal connection relative to some specific distance connecting the topological singularities of $u$ and associated to a measurable weight function. This result generalizes a previous result of F. Bethuel, H. Brezis et J.M. Coron.


Résumé. Nous calculons explicitement une énergie de type relaxée pour des applications $u: \Omega \subset \mathbb{R}^{3} \rightarrow S^{2}$. La formule explicite fait intervenir la longueur d'une connexion minimale relative à une certaine distance, connectant les singularités topologiques de $u$ et associée à une fonction de poids mesurable. Ce résultat généralise un résultat antérieur de F. Bethuel, H. Brezis and J.M. Coron.

MSC: 49D20; 49F99

## 1 Introduction and Main Results

Let $\Omega$ be a smooth bounded and connected open set of $\mathbb{R}^{3}$ and let $w: \Omega \rightarrow \mathbb{R}$ be a measurable function such that

$$
\begin{equation*}
0<\lambda \leq w \leq \Lambda \quad \text { a.e. } \operatorname{in} \Omega \tag{1.1}
\end{equation*}
$$

for some constant $\lambda$ and $\Lambda$. We set $H_{g}^{1}\left(\Omega, S^{2}\right)=\left\{u \in H^{1}\left(\Omega, S^{2}\right), u=g\right.$ on $\left.\partial \Omega\right\}$, where $g: \partial \Omega \rightarrow S^{2}$ is a given smooth boundary data such that $\operatorname{deg}(g)=0$. Our main goal in this paper is to obtain an explicit formula for the relaxed functional
$E_{w}(u)=\operatorname{Inf}\left\{\liminf _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x, u_{n} \in H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega}), u_{n} \rightharpoonup u\right.$ weakly in $\left.H^{1}\right\}$,
defined for $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$. By a result of F . Bethuel (see $\left.[1]\right), H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega})$ is sequentially dense for the weak topology in $H_{g}^{1}\left(\Omega, S^{2}\right)$ and then the functional $E_{w}$ is well defined.

In [4], F. Bethuel, H. Brezis and J.M. Coron have proved that for $w \equiv 1$,

$$
E_{1}(u)=\int_{\Omega}|\nabla u(x)|^{2} d x+8 \pi L(u)
$$

where $L(u)$ denotes the length of a minimal connection relative to the Euclidean geodesic distance $d_{\Omega}$ in $\bar{\Omega}$ connecting the singularities of $u$ (see also M. Giaquinta, G. Modica, J. Souček [13]). If $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ is smooth on $\bar{\Omega}$ except at a finite number of points in $\Omega$, the length of a minimal connection relative to $d_{\Omega}$ connecting the singularities of $u$ is given by

$$
L(u)=\operatorname{Min}_{\sigma \in \mathcal{S}_{K}} \sum_{i=1}^{K} d_{\Omega}\left(P_{i}, N_{\sigma(i)}\right)
$$

where $\left(P_{1}, \ldots, P_{K}\right)$ and $\left(N_{1}, \ldots, N_{K}\right)$ are respectively the singularities of positive and negative degree counted according to their multiplicity (since $\operatorname{deg}(g)=0$, the number of positive singularities is equal to the number of negative ones) and $\mathcal{S}_{K}$ denotes the set of all permutations of $K$ indices. For the definition of $L(u)$ when $u$ is arbitrary in $H_{g}^{1}\left(\Omega, S^{2}\right)$, we refer to (1.6)-(1.7) below. The notion of length of a minimal connection between singularities has its origin in [9]. We also refer to the results of J. Bourgain, H. Brezis, P. Mironescu [5] and H. Brezis, P. Mironescu, A.C. Ponce [10] for similar problems involving $S^{1}$-valued maps.

For $u \in H^{1}\left(\Omega, S^{2}\right)$, the vector field $D(u)$ first introduced in [9] and defined by

$$
\begin{equation*}
D(u)=\left(u \cdot \frac{\partial u}{\partial x_{2}} \wedge \frac{\partial u}{\partial x_{3}}, u \cdot \frac{\partial u}{\partial x_{3}} \wedge \frac{\partial u}{\partial x_{1}}, u \cdot \frac{\partial u}{\partial x_{1}} \wedge \frac{\partial u}{\partial x_{2}}\right) \tag{1.2}
\end{equation*}
$$

plays a crucial role. Indeed, if $u$ is smooth except at a finite number of points $\left(P_{i}, N_{i}\right)_{i=1}^{K}$ in $\Omega$, then (see [9], Appendix B)

$$
\begin{equation*}
\operatorname{div} D(u)=4 \pi \sum_{i=1}^{K}\left(\delta_{P_{i}}-\delta_{N_{i}}\right) \quad \text { in } \mathcal{D}^{\prime}(\Omega) \tag{1.3}
\end{equation*}
$$

and if in addition $u_{\mid \partial \Omega}=g$, we have (since $\operatorname{deg}(g)=0$, see [9], Section IV)

$$
\begin{equation*}
L(u)=\operatorname{Sup}\left\{\sum_{i=1}^{K}\left(\zeta\left(P_{i}\right)-\zeta\left(N_{i}\right)\right)\right\} \tag{1.4}
\end{equation*}
$$

where the supremum is taken over all functions $\zeta: \bar{\Omega} \rightarrow \mathbb{R}$ which are 1-Lipschitz with respect to distance $d_{\Omega}$ i.e., $|\zeta(x)-\zeta(y)| \leq d_{\Omega}(x, y)$. Note that for any real Lipschitz function $\zeta$,

$$
\begin{equation*}
\sum_{i=1}^{K} \zeta\left(P_{i}\right)-\zeta\left(N_{i}\right)=\frac{1}{4 \pi} \int_{\Omega} \operatorname{div} D(u) \zeta=-\frac{1}{4 \pi} \int_{\Omega} D(u) \cdot \nabla \zeta+\frac{1}{4 \pi} \int_{\partial \Omega}(D(u) \cdot \nu) \zeta \tag{1.5}
\end{equation*}
$$

where $\nu$ denotes the outward normal to $\partial \Omega$. We recall that $D(u) \cdot \nu$ is equal to the $2 \times 2$ Jacobian determinant of $u$ restricted to $\partial \Omega$ and then it only depends on $g$. In view of (1.4) and (1.5), $L(u)$ has been defined in [4] for $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ by

$$
\begin{equation*}
L(u)=\frac{1}{4 \pi} \operatorname{Sup}\left\{\langle T(u), \zeta\rangle, \zeta: \bar{\Omega} \rightarrow \mathbb{R} \text { 1-Lipschitz with respect to } d_{\Omega}\right\} \tag{1.6}
\end{equation*}
$$

where $T(u) \in \mathcal{D}^{\prime}(\Omega)$ denotes the distribution defined by its action on real Lipschitz functions through the formula:

$$
\begin{equation*}
\langle T(u), \zeta\rangle=\int_{\Omega} D(u) \cdot \nabla \zeta-\int_{\partial \Omega}(D(u) \cdot \nu) \zeta \tag{1.7}
\end{equation*}
$$

In a previous paper [14], we have studied the following variational problem: given two distinct points $P$ and $N$ in $\Omega$,

$$
E_{w}(P, N)=\operatorname{Inf}\left\{\int_{\Omega}|\nabla v(x)|^{2} w(x) d x, v \in \mathcal{E}(P, N)\right\}
$$

where
$\mathcal{E}(P, N)=\left\{v \in H^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega} \backslash\{P, N\}), v=\right.$ const on $\partial \Omega, T(v)=4 \pi\left(\delta_{P}-\delta_{N}\right)$ in $\left.\mathcal{D}^{\prime}(\Omega)\right\}$.
In the case $w \equiv 1$, H. Brezis, J.M. Coron and E. Lieb have shown that (see [9])

$$
E_{1}(P, N)=8 \pi d_{\Omega}(P, N)
$$

For an arbitrary function $w$, we have proved (see [14]) that $E_{w}(\cdot, \cdot)$ defines a distance function satisfying

$$
\begin{equation*}
8 \pi \lambda d_{\Omega}(\cdot, \cdot) \leq E_{w}(\cdot, \cdot) \leq 8 \pi \Lambda d_{\Omega}(\cdot, \cdot) \tag{1.8}
\end{equation*}
$$

From (1.8), we infer that $E_{w}$ extends to $\bar{\Omega} \times \bar{\Omega}$ into a distance on $\bar{\Omega}$. In what follows, we set for $x, y \in \bar{\Omega}$,

$$
d_{w}(x, y)=\frac{1}{8 \pi} E_{w}(x, y)
$$

When $w$ is continuous, we also have shown that the distance $d_{w}$ can be characterized in the following way: for any $x, y \in \bar{\Omega}$,

$$
d_{w}(x, y)=\operatorname{Min} \int_{0}^{1} w(\gamma(t))|\dot{\gamma}(t)| d t
$$

where the minimum is taken over all Lipschitz curve $\gamma:[0,1] \rightarrow \bar{\Omega}$ verifying $\gamma(0)=x$ and $\gamma(1)=y$. For an arbitrary measurable function $w$, the previous formula is meaningless since $w$ is not well defined on curves but a similar characterization of $d_{w}$ actually holds. We refer to [14] for more details. We also recall the general result in [14]:

Theorem 1.1. Let $\left(P_{i}\right)_{i=1}^{K}$ and $\left(N_{i}\right)_{i=1}^{K}$ be two lists of points in $\Omega$ and consider

$$
\begin{aligned}
\mathcal{E}\left(\left(P_{i}, N_{i}\right)_{i=1}^{K}\right)=\{ & v \in H^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\left(P_{i}, N_{i}\right)_{i=1}^{K}\right\}\right) \\
& \left.v=\mathrm{const} \text { on } \partial \Omega \text { and } T(v)=4 \pi \sum_{i=1}^{K} \delta_{P_{i}}-\delta_{N_{i}} \text { in } \mathcal{D}^{\prime}(\Omega)\right\} .
\end{aligned}
$$

Then we have

$$
\operatorname{Inf}\left\{\int_{\Omega}|\nabla v(x)|^{2} w(x) d x, v \in \mathcal{E}\left(\left(P_{i}, N_{i}\right)_{i=1}^{K}\right)\right\}=8 \pi L_{w}
$$

where $L_{w}$ is the length of a minimal connection relative to distance $d_{w}$ connecting the points $\left(P_{i}\right)$ and $\left(N_{i}\right)$ i.e.,

$$
L_{w}=\operatorname{Min}_{\sigma \in \mathcal{S}_{K}} \sum_{i=1}^{K} d_{w}\left(P_{i}, N_{\sigma(i)}\right)
$$

By analogy with the case $w \equiv 1$, we define for $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$,

$$
L_{w}(u)=\frac{1}{4 \pi} \operatorname{Sup}\left\{\langle T(u), \zeta\rangle, \zeta: \bar{\Omega} \rightarrow \mathbb{R} 1 \text {-Lipschitz with respect to } d_{w}\right\}
$$

(note that any real function $\zeta$ which is 1 -Lipschitz with respect to $d_{w}$, is a Lipschitz function with respect to $d_{\Omega}$ since $d_{w}$ is strongly equivalent to $d_{\Omega}$ and then $\langle T(u), \zeta\rangle$ is well defined). When $u$ is smooth except at a finite number of points $\left(P_{i}, N_{i}\right)_{i=1}^{K}$ in $\Omega$, it follows as in [9] that $L_{w}(u)$ is equal to the length of a minimal connection relative to distance $d_{w}$ connecting the points $\left(P_{i}\right)$ and $\left(N_{i}\right)$. Our main result is the following.

Theorem 1.2. For any $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$, we have

$$
E_{w}(u)=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u)
$$

The proof of Theorem 1.2 is presented in Section 3 and is based on a method similar to the one used in [4] and on a Dipole Removing Technique exposed in the next section. This technique is mostly inspired from [1] but involves some tools developed in [14] in order to treat the problem for a non smooth function $w$.

In Section 4, we prove a stability property of $E_{w}$. More precisely, we give some conditions on a sequence $\left(w_{n}\right)_{n \in \mathbb{N}}$ under which one can conclude that the sequence of functionals $\left(E_{w_{n}}\right)_{n \in \mathbb{N}}$ converges pointwise to $E_{w}$ on $H_{g}^{1}\left(\Omega, S^{2}\right)$. The results are obtained using previous ones in [14]. In Section 5, we present similar results for a relaxed type functional in which we do not prescribed any boundary data.

Throughout the paper, a sequence of smooth mollifiers means any sequence $\left(\rho_{n}\right)_{n \in \mathbb{N}}$ satisfying

$$
\rho_{n} \in \mathcal{C}^{\infty}\left(\mathbb{R}^{3}, \mathbb{R}\right), \quad \text { Supp } \rho_{n} \subset B_{1 / n}, \quad \int_{\mathbb{R}^{3}} \rho_{n}=1, \quad \rho_{n} \geq 0 \text { on } \mathbb{R}^{3}
$$

## 2 The Dipole Removing Technique

In this section, we first give a technical result which will be used for the dipole removing technique in Section 2.2.

### 2.1 Preliminaries

Let $\alpha$ and $\beta$ be two distinct points in $\Omega$. We denote by $p_{\alpha, \beta}(\xi)$ the projection of $\xi \in \mathbb{R}^{3}$ on the straight line passing by $\alpha$ and $\beta$ and $r_{\alpha, \beta}(\xi)=\operatorname{dist}(x,[\alpha, \beta])$, where "dist" denotes the Euclidean distance in $\mathbb{R}^{3}$. For $m \in \mathbb{N}^{*}$, we set

$$
a_{m}^{\alpha, \beta}=\frac{|\alpha-\beta|}{m} \quad \text { and } \quad s_{j}^{\alpha, \beta}=j a_{m}^{\alpha, \beta} \quad \text { for } j=0, \ldots, m
$$

For $\xi \in \mathbb{R}^{3}$ such that $p_{\alpha, \beta}(\xi) \in[\alpha, \beta]$, we define

$$
h_{m}^{\alpha, \beta}(\xi)=\min _{0 \leq j \leq m}| | p_{\alpha, \beta}(\xi)-\alpha\left|-s_{j}^{\alpha, \beta}\right|
$$

and we set

$$
\Theta_{m}([\alpha, \beta])=\left\{\xi \in \mathbb{R}^{3}, p_{\alpha, \beta}(\xi) \in[\alpha, \beta] \text { and } r_{\alpha, \beta}(\xi) \leq a_{m}^{\alpha, \beta} h_{m}^{\alpha, \beta}(\xi)\right\} .
$$

For two points $x$ and $y$ in $\Omega$, we consider the class $\mathcal{Q}(x, y)$ of all finite collections of segments $\mathcal{F}=\left(\left[\alpha_{k}, \beta_{k}\right]\right]_{k=1}^{n(\mathcal{F})}$ such that $\beta_{k}=\alpha_{k+1}, \alpha_{1}=x, \beta_{n(\mathcal{F})}=y,\left[\alpha_{k}, \beta_{k}\right] \subset \Omega$ and $\alpha_{k} \neq \beta_{k}$. We define the "length" of an element $\mathcal{F} \in \mathcal{Q}(x, y)$ by

$$
\bar{\ell}_{w}(\mathcal{F})=\sum_{k=1}^{n(\mathcal{F})} \liminf _{m \rightarrow+\infty} \frac{1}{\pi} \int_{\Theta_{m}\left(\left[\alpha_{k}, \beta_{k}\right]\right) \cap \Omega} \varepsilon_{\alpha_{k}, \beta_{k}}^{m}(\xi) w(\xi) d \xi
$$

with

$$
\varepsilon_{\alpha_{k}, \beta_{k}}^{m}(\xi)=\frac{\left(h_{m}^{\alpha_{k}, \beta_{k}}(\xi)\right)^{2}\left(a_{m}^{\alpha_{k}, \beta_{k}}\right)^{4}}{\left(\left(h_{m}^{\alpha_{k}, \beta_{k}}(\xi)\right)^{2}\left(a_{m}^{\alpha_{k}, \beta_{k}}\right)^{4}+r_{\alpha_{k}, \beta_{k}}^{2}(\xi)\right)^{2}} .
$$

Lemma 2.1. Let $\mathbb{P}$ be a finite collection of distinct points in $\Omega$ or $\mathbb{P}=\emptyset$. For any distinct points $x_{0}, y_{0}$ in $\Omega \backslash \mathbb{P}$ and $\delta>0$, there exists $\mathcal{F}_{\delta}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right) \in \mathcal{Q}\left(x_{0}, y_{0}\right)$ such that $\left(\mathbb{P} \cup\left\{y_{0}\right\}\right) \cap\left(\cup_{k=1}^{n-1}\left[\alpha_{k}, \beta_{k}\right] \cup\left[\alpha_{n}, \beta_{n}[)=\emptyset\right.\right.$ and

$$
\bar{\ell}_{w}(\mathcal{F}) \leq d_{w}\left(x_{0}, y_{0}\right)+\delta
$$

Proof. Step 1. Assume that $w$ is smooth on $\Omega$. We are going to prove that for every element $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right) \in \mathcal{Q}(x, y)$, we have

$$
\bar{\ell}_{w}(\mathcal{F})=\int_{\bigcup_{k=1}^{n}\left[\alpha_{k}, \beta_{k}\right]} w(s) d s .
$$

It suffices to prove that for any distinct points $\alpha, \beta \in \Omega$,

$$
\begin{equation*}
\lim _{m \rightarrow+\infty} \frac{1}{\pi} \int_{\Theta_{m}([\alpha, \beta]) \cap \Omega} \varepsilon_{k}^{m}(\xi) w(\xi) d \xi=\int_{[\alpha, \beta]} w(s) d s \tag{2.1}
\end{equation*}
$$

Without loss of generality, we may assume that $[\alpha, \beta]=\{(0,0)\} \times[0, R]$ and we drop the indices $\alpha$ and $\beta$ for simplicity. We set for $j=0, \ldots, m-1$,

$$
C_{m}^{j+}=\left\{\xi=\left(\xi_{1}, \xi_{2}, \xi_{3}\right) \in \Theta_{m}([\alpha, \beta]), \xi_{3} \in\left[s_{j}, s_{j}+\frac{a_{m}}{2}\right]\right\},
$$

and for $j=1, \ldots, m$,

$$
C_{m}^{j-}=\left\{\xi=\left(\xi_{1}, \xi_{2}, \xi_{3}\right) \in \Theta_{m}([\alpha, \beta]), \xi_{3} \in\left[s_{j}-\frac{a_{m}}{2}, s_{j}\right]\right\} .
$$

For $\xi \in C_{m}^{j+} \cup C_{m}^{j-}$, we have $h_{m}(\xi)=\left|\xi_{3}-s_{j}\right|$ and we get that for $m$ large enough,

$$
\begin{equation*}
\int_{\Theta_{m}([\alpha, \beta]) \cap \Omega} \varepsilon_{k}^{m}(\xi) w(\xi) d \xi=\sum_{j=0}^{m-1} I_{m}^{j+}+\sum_{j=1}^{m} I_{m}^{j-} \tag{2.2}
\end{equation*}
$$

with

$$
\begin{aligned}
I_{m}^{j+} & =\int_{C_{m}^{j+}} \frac{\left|\xi_{3}-s_{j}\right|^{2} a_{m}^{4} w(\xi)}{\left(\left|\xi_{3}-s_{j}\right|^{2} a_{m}^{4}+r^{2}(\xi)\right)^{2}} d \xi \quad \text { for } j=0, \ldots, m-1 \\
I_{m}^{j-} & =\int_{C_{m}^{j-}} \frac{\left|\xi_{3}-s_{j}\right|^{2} a_{m}^{4} w(\xi)}{\left(\left|\xi_{3}-s_{j}\right|^{2} a_{m}^{4}+r^{2}(\xi)\right)^{2}} d \xi \quad \text { for } j=1, \ldots, m
\end{aligned}
$$

Using the change of variable $z_{1}=\frac{\xi_{1}}{\left|\xi_{3}-s_{j}\right|}, z_{2}=\frac{\xi_{2}}{\left|\xi_{3}-s_{j}\right|}$ and $z_{3}=\xi_{3}$, we derive that

$$
\begin{aligned}
I_{m}^{j+} & =\int_{s_{j}}^{s_{j}+\frac{a_{m}}{2}}\left(\int_{B_{a_{m}}(0)} \frac{a_{m}^{4} w\left(\left|z_{3}-s_{j}\right| z_{1},\left|z_{3}-s_{j}\right| z_{2}, z_{3}\right)}{\left(a_{m}^{4}+z_{1}^{2}+z_{2}^{2}\right)^{2}} d z_{1} d z_{2}\right) d z_{3} \\
& =\int_{s_{j}}^{s_{j}+\frac{a_{m}}{2}}\left(w\left(0,0, z_{3}\right)+\mathcal{O}\left(a_{m}\right)\right)\left(\int_{B_{a_{m}}(0)} \frac{a_{m}^{4}}{\left(a_{m}^{4}+z_{1}^{2}+z_{2}^{2}\right)^{2}} d z_{1} d z_{2}\right) d z_{3} \\
& =\pi \int_{s_{j}}^{s_{j}+\frac{a_{m}}{2}} w\left(0,0, z_{3}\right) d z_{3}+\mathcal{O}\left(a_{m}^{2}\right)
\end{aligned}
$$

By similar computations we get that

$$
I_{m}^{j-}=\pi \int_{s_{j}-\frac{a_{m}}{2}}^{s_{j}} w\left(0,0, z_{3}\right) d z_{3}+\mathcal{O}\left(a_{m}^{2}\right)
$$

Combining this equalities with (2.2), we obtain that

$$
\int_{\Theta_{m}([\alpha, \beta]) \cap \Omega} \varepsilon_{k}^{m}(\xi) w(\xi) d \xi=\pi \int_{0}^{R} w\left(0,0, z_{3}\right) d z_{3}+\mathcal{O}\left(a_{m}\right)
$$

which ends the proof of (2.1).
Step 2. We fix two distinct points $x_{0}, y_{0} \in \Omega \backslash \mathbb{P}$. For any points $x, y$ in $\Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$, let $\mathcal{Q}^{\prime}(x, y)$ be the class of elements $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right) \in \mathcal{Q}(x, y)$ such that

$$
\cup_{k=1}^{n}\left[\alpha_{k}, \beta_{k}\right] \subset \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right) .
$$

We consider the function $\mathcal{D}_{w}: \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right) \times \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right) \rightarrow \mathbb{R}_{+}$defined by

$$
\mathcal{D}_{w}(x, y)=\inf _{\mathcal{F} \in \mathcal{Q}^{\prime}(x, y)} \bar{\ell}(\mathcal{F})
$$

We are going to show that $\mathcal{D}_{w}$ defines a distance function which can be extended to $\bar{\Omega} \times \bar{\Omega}$. Let $x, y \in \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$ and let $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right)$ be an element of $\mathcal{Q}^{\prime}(x, y)$. Assumption (1.1) and similar computations to those in Step 1 lead to

$$
\lambda \sum_{k=1}^{n}\left|\alpha_{k}-\beta_{k}\right| \leq \bar{\ell}_{w}(\mathcal{F}) \leq \Lambda \sum_{k=1}^{n}\left|\alpha_{k}-\beta_{k}\right|
$$

Taking the infimum over all $\mathcal{F} \in \mathcal{Q}^{\prime}(x, y)$, we infer that

$$
\begin{equation*}
\lambda d_{\Omega}(x, y) \leq \mathcal{D}_{w}(x, y) \leq \Lambda d_{\Omega}(x, y) \tag{2.3}
\end{equation*}
$$

From (2.3), we deduce that $\mathcal{D}_{w}(x, y)=0$ if and only if $x=y$. Let us now prove that $\mathcal{D}_{w}$ is symmetric. Let $x, y \in \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$ and $\delta>0$ arbitrary small. By definition, we can find $\mathcal{F}_{\delta}=\left(\left[\alpha_{1}, \beta_{2}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right)$ in $\mathcal{Q}^{\prime}(x, y)$ satisfying

$$
\bar{\ell}_{w}\left(\mathcal{F}_{\delta}\right) \leq \mathcal{D}_{w}(x, y)+\delta
$$

Then for $\mathcal{F}_{\delta}^{\prime}=\left(\left[\beta_{n}, \alpha_{n}\right], \ldots,\left[\beta_{1}, \alpha_{1}\right]\right) \in \mathcal{Q}^{\prime}(y, x)$, we have

$$
\mathcal{D}_{w}(y, x) \leq \bar{\ell}_{w}\left(\mathcal{F}_{\delta}^{\prime}\right)=\bar{\ell}_{w}\left(\mathcal{F}_{\delta}\right) \leq \mathcal{D}_{w}(x, y)+\delta
$$

Since $\delta$ is arbitrary, we obtain $\mathcal{D}_{w}(y, x) \leq \mathcal{D}_{w}(x, y)$ and we conclude that $\mathcal{D}_{w}(y, x)=\mathcal{D}_{w}(x, y)$ inverting the roles of $x$ and $y$. The triangle inequality is immediate since the juxtaposition of $\mathcal{F}_{1} \in \mathcal{Q}^{\prime}(x, z)$ with $\mathcal{F}_{2} \in \mathcal{Q}^{\prime}(z, y)$ is an element of $\mathcal{Q}^{\prime}(x, y)$. Hence $\mathcal{D}_{w}$ defines a distance on $\Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$ verifying (2.3). Therefore distance $\mathcal{D}_{w}$ extends uniquely to $\bar{\Omega} \times \bar{\Omega}$ into a distance function that we still denote by $\mathcal{D}_{w}$. By continuity, $\mathcal{D}_{w}$ satisfies (2.3) for any $x, y \in \bar{\Omega}$.
Step 3. We consider the function $\zeta: \bar{\Omega} \rightarrow \mathbb{R}$ defined by

$$
\zeta(x)=\mathcal{D}_{w}\left(x, x_{0}\right)
$$

Note that function $\zeta$ is 1 -Lipschitz with respect to distance $\mathcal{D}_{w}$ and therefore $\Lambda$-Lipschitz with respect to the Euclidean geodesic distance on $\bar{\Omega}$ by (2.3). We fix an arbitrary point $z_{0} \in \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$ and some $R>0$ such that $B_{3 R}\left(z_{0}\right) \subset \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$. Let $\left(\rho_{n}\right)_{n \in \mathbb{N}}$ be a sequence of smooth mollifiers. For $n>1 / R$, we consider the smooth function $\zeta_{n}=\rho_{n} * \zeta$ : $B_{R}\left(z_{0}\right) \rightarrow \mathbb{R}$. We write

$$
\zeta_{n}(x)=\int_{B_{1 / n}} \rho_{n}(-z) \zeta(x+z) d z
$$

and therefore for all $x, y \in B_{R}\left(z_{0}\right)$,

$$
\begin{aligned}
\left|\zeta_{n}(x)-\zeta_{n}(y)\right| & \leq \int_{B_{1 / n}} \rho_{n}(-z)|\zeta(x+z)-\zeta(y+z)| d z \\
& \leq \int_{B_{1 / n}} \rho_{n}(-z) \mathcal{D}_{w}(x+z, y+z) d z \\
& \leq \int_{B_{1 / n}} \rho_{n}(-z) \bar{\ell}_{w}([x+z, y+z]) d z
\end{aligned}
$$

We remark that $\Theta_{m}([x+z, y+z])=z+\Theta_{m}([x, y])$. For $m$ large enough $z+\Theta_{m}([x, y]) \subset$ $B_{3 R}\left(z_{0}\right)$ and then for any vector $\xi \in \Theta_{m}([x, y])$, we have $\varepsilon_{x+z, y+z}^{m}(\xi+z)=\varepsilon_{x, y}^{m}(\xi)$. Hence we obtain for all $z \in B_{1 / n}(0)$ and $m$ sufficiently large,

$$
\bar{\ell}_{w}([x+z, y+z])=\liminf _{m \rightarrow+\infty} \frac{1}{\pi} \int_{\Theta_{m}([x, y])} \varepsilon_{x, y}^{m}(\xi) w(\xi+z) d \xi
$$

Using Fatou's lemma, we get that

$$
\begin{aligned}
\left|\zeta_{n}(x)-\zeta_{n}(y)\right| & \leq \int_{B_{1 / n}} \rho_{n}(-z)\left(\liminf _{m \rightarrow+\infty} \frac{1}{\pi} \int_{\Theta_{m}([x, y])} \varepsilon_{x, y}^{m}(\xi) w(\xi+z) d \xi\right) d z \\
& \leq \liminf _{m \rightarrow+\infty} \frac{1}{\pi} \int_{B_{1 / n}} \int_{\Theta_{m}([x, y])} \rho_{n}(-z) \varepsilon_{x, y}^{m}(\xi) w(\xi+z) d \xi d z
\end{aligned}
$$

For each $m \in \mathbb{N}$ sufficiently large we have

$$
\frac{1}{\pi} \int_{B_{1 / n}} \int_{\Theta_{m}([x, y])} \rho_{n}(-z) \varepsilon_{x, y}^{m}(\xi) w(\xi+z) d \xi d z=\frac{1}{\pi} \int_{\Theta_{m}([x, y])} \varepsilon_{x, y}^{m}(\xi) \rho_{n} * w(\xi) d \xi
$$

and since $\rho_{n} * w$ is smooth, we obtain as in Step 1,

$$
\frac{1}{\pi} \int_{\Theta_{m}([x, y])} \varepsilon_{x, y}^{m}(\xi) \rho_{n} * w(\xi) d \xi \rightarrow \int_{[x, y]} \rho_{n} * w(s) d s \quad \text { as } m \rightarrow+\infty
$$

Thus for each $x, y \in B_{R}\left(z_{0}\right)$ we have

$$
\left|\zeta_{n}(x)-\zeta_{n}(y)\right| \leq \int_{[x, y]} \rho_{n} * w(s) d s
$$

Then for $x \in B_{R}\left(z_{0}\right), h \in S^{2}$ fixed and $\delta>0$ small, we infer that

$$
\frac{\left|\zeta_{n}(x+\delta h)-\zeta_{n}(x)\right|}{\delta} \leq \frac{1}{\delta} \int_{[x, x+\delta h]} \rho_{n} * w(s) d s \underset{\delta \rightarrow 0^{+}}{\rightarrow} \rho_{n} * w(x)
$$

and we conclude, letting $\delta \rightarrow 0$, that $\left|\nabla \zeta_{n}(x) \cdot h\right| \leq \rho_{n} * w(x)$ for each $x \in B_{R}\left(z_{0}\right)$ and $h \in S^{2}$ which implies that $\left|\nabla \zeta_{n}\right| \leq \rho_{n} * w$ on $B_{R}\left(z_{0}\right)$. Since $\nabla \zeta_{n} \rightarrow \nabla \zeta$ and $\rho_{n} * w \rightarrow w$ a.e. on $B_{R}\left(z_{0}\right)$ as $n \rightarrow+\infty$, we deduce that $|\nabla \zeta| \leq w$ a.e. on $B_{R}\left(z_{0}\right)$. Since $z_{0}$ is arbitrary in $\Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$, we derive

$$
|\nabla \zeta| \leq w \quad \text { a.e. on } \Omega \text {. }
$$

By Proposition 2.3. in [14], it follows that $|\zeta(x)-\zeta(y)| \leq d_{w}(x, y)$ for any $x, y \in \bar{\Omega}$ and in particular, we obtain choosing $y=x_{0}$,

$$
\mathcal{D}_{w}\left(x, x_{0}\right) \leq d_{w}\left(x, x_{0}\right) \quad \text { for all } x \in \bar{\Omega} .
$$

Step 4. End of the Proof. Let $\delta>0$ be given. We choose some $\tilde{y}_{0} \in \Omega \backslash\left(\mathbb{P} \cup\left\{y_{0}\right\}\right)$ such that $\left[\tilde{y}_{0}, y_{0}\right] \subset \Omega \backslash \mathbb{P}$ and $\left|\tilde{y}_{0}-y_{0}\right| \leq \frac{\delta}{3 \Lambda}$. By the previous step, we can find an element $\mathcal{F}^{\prime}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right) \in \mathcal{Q}^{\prime}\left(x_{0}, \tilde{y}_{0}\right)$ verifying

$$
\bar{\ell}_{w}\left(\mathcal{F}^{\prime}\right) \leq d_{w}\left(x_{0}, \tilde{y}_{0}\right)+\frac{\delta}{3} .
$$

Then we consider $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right],\left[\tilde{y}_{0}, y_{0}\right]\right) \in \mathcal{Q}\left(x_{0}, y_{0}\right)$. We have

$$
\begin{aligned}
\bar{\ell}_{w}(\mathcal{F}) \leq \bar{\ell}_{w}\left(\mathcal{F}^{\prime}\right)+\Lambda\left|\tilde{y}_{0}-y_{0}\right| & \leq d_{w}\left(x_{0}, \tilde{y}_{0}\right)+\frac{2 \delta}{3} \\
& \leq d_{w}\left(x_{0}, y_{0}\right)+d_{w}\left(y_{0}, \tilde{y}_{0}\right)+\frac{2 \delta}{3} \\
& \leq d_{w}\left(x_{0}, y_{0}\right)+\delta
\end{aligned}
$$

and then $\mathcal{F}$ satisfies the requirement.

### 2.2 The Dipole Removing Technique

We first present the dipole removing technique for a simple dipole. We then treat the case of several point singularities.

Lemma 2.2. Let $P$ and $N$ be two distinct points in $\Omega$ and consider $u \in H^{1}\left(\Omega, S^{2}\right) \cap$ $\mathcal{C}^{1}(\bar{\Omega} \backslash\{P, N\})$ with $\operatorname{deg}(u, P)=+1$ and $\operatorname{deg}(u, N)=-1$. Let $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right)$ be an element of $\mathcal{Q}(P, N)$ such that $N \notin \cup_{k=1}^{n-1}\left[\alpha_{k}, \beta_{k}\right] \cup\left[\alpha_{n}, \beta_{n}[\right.$. Then for any $\delta>0$ small enough, there exists a map $u_{\delta} \in \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ such that:

$$
\int_{\Omega}\left|\nabla u_{\delta}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \bar{\ell}_{w}(\mathcal{F})+\delta
$$

and $u_{\delta}$ coincides with $u$ outside a $\delta$-neighborhood of $\cup_{k=1}^{n}\left[\alpha_{k}, \beta_{k}\right]$ included in $\Omega$.
Proof. Let $\mathcal{F}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{n}, \beta_{n}\right]\right) \in \mathcal{Q}(P, N)$ such that $N \notin \cup_{k=1}^{n-1}\left[\alpha_{k}, \beta_{k}\right] \cup\left[\alpha_{n}, \beta_{n}[\right.$ and fix some $\delta>0$ small. We proceed in several steps.
Step 1. We consider a small $0<r_{0}<\delta$ verifying $B_{r_{0}}\left(\alpha_{1}\right) \subset \Omega \backslash\{N\}$. By Lemma A. 1 in [1], we can find $v \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\alpha_{1}, N\right\}, S^{2}\right) \cap H^{1}(\Omega)$ (recall that $\left.\alpha_{1}=P\right)$ satisfying

$$
v(x)= \begin{cases}u(x) & \text { on } \Omega \backslash B_{r_{0}}\left(\alpha_{1}\right),  \tag{2.4}\\ R\left(\frac{x-\alpha_{1}}{\left|x-\alpha_{1}\right|}\right) & \text { on } B_{r_{0}}\left(\alpha_{1}\right),\end{cases}
$$

for some rotation $R$ and

$$
\begin{equation*}
\int_{\Omega}|\nabla v(x)|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+\delta . \tag{2.5}
\end{equation*}
$$

Let $W=\left\{x \in \mathbb{R}^{3}\right.$, $\left.\operatorname{dist}\left(x,\left[\alpha_{1}, \beta_{1}\right]\right)<\delta\right\}$. For $\delta$ small enough, we have $\bar{W} \subset \Omega \backslash\{N\}$. We set $d=\left|\alpha_{1}-\beta_{1}\right|$. We choose normal coordinates such that $\alpha_{1}=(0,0,0)$ and $\beta_{1}=(0,0, d)$. Let $0<r<\frac{r_{0}}{2}$. Since $v$ is smooth on $W \backslash B_{r_{0}}\left(\alpha_{1}\right)$, we can find a constant $\sigma(r)$ such that $|\nabla v| \leq \sigma(r)$ on $W \backslash B_{r_{0}}\left(\alpha_{1}\right)$. For $m \in \mathbb{N}^{*}$, we consider

$$
K_{m}=\left[-\frac{a_{m}^{\alpha_{1}, \beta_{1}}}{2}, \frac{a_{m}^{\alpha_{1}, \beta_{1}}}{2}\right]^{2} \times\left[-\frac{a_{m}^{\alpha_{1}, \beta_{1}}}{2}, d+\frac{a_{m}^{\alpha_{1}, \beta_{1}}}{2}\right] .
$$

For $m$ large enough, we have $\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right) \subset K_{m} \subset W$. As in [1], we are going to construct in the next step a map $v_{1} \in \mathcal{C}^{1}\left(\bar{W} \backslash\left\{\beta_{1}\right\}, S^{2}\right) \cap H^{1}(W)$ verifying $v_{1}=v$ in a neighborhood of $\partial W$ and $\operatorname{deg}\left(v_{1}, \beta_{1}\right)=+1$. For simplicity, we drop the indices $\alpha_{1}$ and $\beta_{1}$.
Step 2. We divide $K_{m}$ in $m+1$ cubes $Q_{m}^{j}$ defined by

$$
Q_{m}^{j}=\left[-\frac{a_{m}}{2}, \frac{a_{m}}{2}\right]^{2} \times\left[\left(j-\frac{1}{2}\right) a_{m},\left(j+\frac{1}{2}\right) a_{m}\right] \quad \text { for } j=0, \ldots, m .
$$

Arguing as in [1], we infer from (2.4) that

$$
\begin{equation*}
\sum_{j=0}^{m} \int_{\partial Q_{m}^{j}}|\nabla v|^{2} \leq C\left(\frac{r}{a_{m}}+m \sigma(r)^{2} a_{m}^{2}\right) . \tag{2.6}
\end{equation*}
$$

We are going to make use of a map $\omega_{m}: B_{a_{m}}^{2}(0) \subset \mathbb{R}^{2} \rightarrow S^{2}$ defined by

$$
\omega_{m}\left(x_{1}, x_{2}\right)=\frac{2 a_{m}^{2}}{a_{m}^{4}+x_{1}^{2}+x_{2}^{2}}\left(x_{1}, x_{2},-a_{m}^{2}\right)+(0,0,1)
$$

( $\omega_{m}$ was first introduced in [8] and we refer to the proof of Lemma 2 in [8] for its main properties). For $j=1, \ldots, m$, we choose an orthonormal direct basis $\left(e_{1}^{j}, e_{2}^{j}, e_{3}^{j}\right)$ of $\mathbb{R}^{3}$ such that

$$
v\left(0,0,(j-1 / 2) a_{m}\right)=(0,0,1) \quad \text { in the basis }\left(e_{1}^{j}, e_{2}^{j}, e_{3}^{j}\right),
$$

and we define the map $v_{1}^{m}: \cup_{j=0}^{m} \partial Q_{m}^{j} \rightarrow S^{2}$ by

1) for $\left(x_{1}, x_{2}, x_{3}\right) \in\left(\cup_{j=0}^{m} \partial Q_{m}^{j}\right) \backslash\left(\cup_{j=1}^{m} B_{a_{m}^{2}}^{2}(0) \times\left\{(j-1 / 2) a_{m}\right\}\right)$,

$$
v_{1}^{m}\left(x_{1}, x_{2}, x_{3}\right)=v\left(x_{1}, x_{2}, x_{3}\right),
$$

2) for $j=1, \ldots, m$ and $\left(x_{1}, x_{2}, x_{3}\right) \in B_{\frac{a_{m}^{2}}{2}}^{2}(0) \times\left\{(j-1 / 2) a_{m}\right\}$,

$$
v_{1}^{m}\left(x_{1}, x_{2}, x_{3}\right)=\omega_{m}\left(\frac{2 x_{1}}{a_{m}}, \frac{2 x_{2}}{a_{m}}\right) \quad \text { in the basis }\left(e_{1}^{j}, e_{2}^{j}, e_{3}^{j}\right),
$$

3) for $j=1, \ldots, m$, for $\left(x_{1}, x_{2}, x_{3}\right) \in\left(B_{a_{m}^{2}}^{2}(0) \backslash B_{\frac{a_{m}^{2}}{2}}^{2}(0)\right) \times\left\{(j-1 / 2) a_{m}\right\}$ and using cylindrical coordinates $\left(x_{1}, x_{2}, x_{3}\right)=(\rho \cos \theta, \rho \sin \theta, z)$,

$$
v_{1}^{m}\left(x_{1}, x_{2}, x_{3}\right)=\left(A_{1} \rho+B_{1}, A_{2} \rho+B_{2}, \sqrt{1-\left(A_{1} \rho+B_{1}\right)^{2}-\left(A_{2} \rho+B_{2}\right)^{2}}\right)
$$

in the basis $\left(e_{1}^{j}, e_{2}^{j}, e_{3}^{j}\right)$, where $A_{1}, A_{2}, B_{1}, B_{2}$ are determined to make $v_{1}^{m}$ continuous. More precisely, if we write $v=v_{1} e_{1}^{j}+v_{2} e_{2}^{j}+v_{3} e_{3}^{j}$ then

$$
\left\{\begin{array}{l}
a_{m}^{2} A_{i}(\theta)+B_{i}(\theta)=v_{i}\left(a_{m}^{2} \cos \theta, a_{m}^{2} \sin \theta,(j-1 / 2) a_{m}\right) \quad i=1,2 \\
\frac{a_{m}^{2}}{2} A_{1}(\theta)+B_{1}(\theta)=\frac{2 a_{m}^{3}}{a_{m}^{4}+a_{m}^{2}} \cos \theta \\
\frac{a_{m}^{2}}{2} A_{2}(\theta)+B_{2}(\theta)=\frac{2 a_{m}^{3}}{a_{m}^{4}+a_{m}^{2}} \sin \theta
\end{array}\right.
$$

The map $v_{1}^{m}$ satisfies by construction $v_{1}^{m}=v$ on $\partial K_{m}$. Moreover, it follows exactly as in the proof of Lemma 2 in [1] that $\operatorname{deg}\left(v_{1}^{m}, \partial Q_{m}^{j}\right)=0$ for $j=0, \ldots, m-1$ and $\operatorname{deg}\left(v_{1}^{m}, \partial Q_{m}^{m}\right)=+1$. Then we extend $v_{1}^{m}$ on each cube $Q_{m}^{j}$ by setting

$$
v_{1}^{m}(x)=v_{1}^{m}\left(\frac{a_{m}\left(x-b_{j}\right)}{2\left\|x-b_{j}\right\|_{\infty}}+b_{j}\right) \quad \text { on } Q_{m}^{j} \text { for } j=0, \ldots, m,
$$

where $b_{j}=\left(0,0, s_{j}\right)$ is the barycenter of $Q_{m}^{j}$ and $\left\|x-b_{j}\right\|_{\infty}=\max \left(\left|x_{1}\right|,\left|x_{2}\right|,\left|x_{3}-s_{j}\right|\right)$. We easily check that $v_{1}^{m} \in H^{1}\left(K_{m}, S^{2}\right), v_{1}^{m}=v$ on $\partial K_{m}, v_{1}^{m}$ is continuous except at the points $b_{j}$ and Lipschitz continuous outside any small neighborhood of the points $b_{j}$. We also get that

$$
\begin{equation*}
\operatorname{deg}\left(v_{1}^{m}, b_{m}\right)=+1 \quad \text { and } \quad \operatorname{deg}\left(v_{1}^{m}, b_{j}\right)=0 \quad \text { for } j=0, \ldots, m-1 . \tag{2.7}
\end{equation*}
$$

We remark that if we set

$$
\begin{gathered}
D_{m}^{j}=B_{\frac{a_{m}^{2}}{2}}^{2}(0) \times\left\{(j-1 / 2) a_{m}\right\} \cup B_{\frac{a_{m}^{2}}{2}}^{2}(0) \times\left\{(j+1 / 2) a_{m}\right\} \quad \text { for } j=1, \ldots, m-1, \\
D_{m}^{0}=B_{\frac{a_{m}^{2}}{2}}^{2}(0) \times\left\{1 / 2 a_{m}\right\} \quad \text { and } \quad D_{m}^{m}=B_{\frac{a_{m}^{2}}{2}}^{2}(0) \times\left\{(m-1 / 2) a_{m}\right\},
\end{gathered}
$$

then we have

$$
\bigcup_{j=0}^{m}\left\{x \in Q_{m}^{j}, \frac{a_{m}\left(x-b_{j}\right)}{2\left\|x-b_{j}\right\|_{\infty}}+b_{j} \in D_{m}^{j} \text { if } x \neq b_{j} \text { or } x=b_{j} \text { otherwise }\right\}=\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)
$$

and if $x \in Q_{m}^{j} \cap \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)$ for some $j \in\{0, \ldots, m\}$ then

$$
\begin{equation*}
h_{m}(x)=\left|x_{3}-s_{j}\right|=\left\|x-b_{j}\right\|_{\infty} \quad \text { and } \quad r(x)=\sqrt{x_{1}^{2}+x_{2}^{2}} \tag{2.8}
\end{equation*}
$$

Some classical computations (see [1] and [8]) lead to, for $j=0, \ldots, m$,

$$
\int_{\left(\partial Q_{m}^{j}\right) \backslash D_{m}^{j}}\left|\nabla v_{1}^{m}\right|^{2} \leq \int_{\partial Q_{m}^{j}}|\nabla v|^{2}+\mathcal{O}\left(a_{m}^{2}\right)
$$

and therefore

$$
\int_{Q_{m}^{j} \backslash \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)}\left|\nabla v_{1}^{m}(x)\right|^{2} w(x) d x \leq C_{1} \Lambda a_{m} \int_{\partial Q_{m}^{j}}|\nabla v|^{2}+C_{2} \Lambda a_{m}^{3}
$$

Adding these inequalities for $j=0, \ldots, m$ and combining with (2.6) we obtain

$$
\begin{equation*}
\int_{K_{m} \backslash \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)}\left|\nabla v_{1}^{m}(x)\right|^{2} w(x) d x \leq C \Lambda\left(r+m \sigma(r)^{2} a_{m}^{3}+a_{m}^{2}\right) \tag{2.9}
\end{equation*}
$$

For $x \in Q_{m}^{j} \cap \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)$ for some $j \in\{0, \ldots, m\}$, we have

$$
v_{1}^{m}(x)= \begin{cases}\omega_{m}\left(\frac{x_{1}}{\left|x_{3}-s_{j}\right|}, \frac{x_{2}}{\left|x_{3}-s_{j}\right|}\right) & \text { in the basis }\left(e_{1}^{j+1}, e_{2}^{j+1}, e_{3}^{j+1}\right) \text { if } x_{3}-s_{j}>0 \\ \omega_{m}\left(\frac{x_{1}}{\left|x_{3}-s_{j}\right|}, \frac{x_{2}}{\left|x_{3}-s_{j}\right|}\right) & \text { in the basis }\left(e_{1}^{j}, e_{2}^{j}, e_{3}^{j}\right) \text { otherwise. }\end{cases}
$$

Following the computations in [6], we infer that

$$
\left|\nabla v_{1}^{m}(x)\right|^{2} \leq \frac{1+C a_{m}^{2}}{\left|x_{3}-s_{j}\right|^{2}}\left|\nabla \omega_{m}\left(\frac{x_{1}}{\left|x_{3}-s_{j}\right|}, \frac{x_{2}}{\left|x_{3}-s_{j}\right|}\right)\right|^{2} \quad \text { in } Q_{m}^{j} \cap \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)
$$

Since we have (see [8])

$$
\left|\nabla \omega_{m}\left(\frac{x_{1}}{\left|x_{3}-s_{j}\right|}, \frac{x_{2}}{\left|x_{3}-s_{j}\right|}\right)\right|^{2}=\frac{8\left|x_{3}-s_{j}\right|^{4} a_{m}^{4}}{\left(\left|x_{3}-s_{j}\right|^{2} a_{m}^{4}+x_{1}^{2}+x_{2}^{2}\right)^{2}}
$$

we derive that

$$
\int_{Q_{m}^{j} \cap \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)}\left|\nabla v_{1}^{m}(x)\right|^{2} w(x) d x \leq \int_{Q_{m}^{j} \cap \Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)} \frac{8\left|x_{3}-s_{j}\right|^{2} a_{m}^{4} w(x)}{\left(\left|x_{3}-s_{j}\right|^{2} a_{m}^{4}+x_{1}^{2}+x_{2}^{2}\right)^{2}} d x+C \Lambda a_{m}^{3}
$$

Summing these inequalities for $j=0, \ldots, m$ and using (2.8) we obtain that

$$
\begin{equation*}
\int_{\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)}\left|\nabla v_{1}^{m}(x)\right|^{2} w(x) d x \leq 8 \int_{\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)} \varepsilon_{\alpha_{1}, \beta_{1}}^{m}(x) w(x) d x+C \Lambda a_{m}^{2} \tag{2.10}
\end{equation*}
$$

Combining (2.9) with (2.10) we conclude that

$$
\int_{K_{m}}\left|\nabla v_{1}^{m}(x)\right|^{2} w(x) d x \leq 8 \int_{\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)} \varepsilon_{\alpha_{1}, \beta_{1}}^{m}(x) w(x) d x+C \Lambda\left(r+m \sigma(r)^{2} a_{m}^{3}+a_{m}^{2}\right) .
$$

Taking the liminf in $m$, we derive that we can find $m_{1} \in \mathbb{N}$ large and $r$ small enough such that

$$
\begin{equation*}
\int_{K_{m_{1}}}\left|\nabla v_{1}^{m_{1}}(x)\right|^{2} w(x) d x \leq 8 \liminf _{m \rightarrow+\infty} \int_{\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)} \varepsilon_{\alpha_{1}, \beta_{1}}^{m}(x) w(x) d x+\delta . \tag{2.11}
\end{equation*}
$$

Since $v_{1}^{m_{1}}=v$ on $\partial K_{m_{1}}$, we may extend $v_{1}^{m_{1}}$ to $W$ by setting $v_{1}^{m_{1}}=v$ on $W \backslash K_{m_{1}}$. Now we recall that $v_{1}^{m_{1}}$ is singular only at the points $b_{j}, j=0, \ldots, m$ (we also recall that $b_{m}=\beta_{1}$ ). From (2.7) and the results in $[1,2,3]$, we infer that exists a map $v_{1} \in \mathcal{C}^{1}\left(\bar{W} \backslash\left\{\beta_{1}\right\}, S^{2}\right) \cap$ $H^{1}(W)$ satisfying $v_{1}=v$ in a neighborhood of $\partial W, \operatorname{deg}\left(v_{1}, \beta_{1}\right)=+1$ and

$$
\begin{equation*}
\int_{W_{1}}\left|\nabla v_{1}(x)\right|^{2} w(x) d x \leq \int_{W_{1}}\left|\nabla v_{1}^{m_{1}}(x)\right|^{2} w(x) d x+\delta . \tag{2.12}
\end{equation*}
$$

Since $v=u$ in a neighborhood of $\partial W$, we may extend $v_{1}$ to $\bar{\Omega}$ by setting $v_{1}=u$ on $\bar{\Omega} \backslash W$. Then we conclude that $v_{1} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\beta_{1}, N\right\}, S^{2}\right) \cap H^{1}(\Omega), \operatorname{deg}\left(v_{1}, \beta_{1}\right)=+1, \operatorname{deg}\left(v_{1}, N\right)=-1$ and by (2.5)-(2.11)-(2.12),

$$
\int_{\Omega}\left|\nabla v_{1}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \liminf _{m \rightarrow+\infty} \int_{\Theta_{m}\left(\left[\alpha_{1}, \beta_{1}\right]\right)} \varepsilon_{\alpha_{1}, \beta_{1}}^{m}(x) w(x) d x+C \delta .
$$

Step 3. Applying Step 1 and Step 2 to $v_{1}$ instead of $u$ and replacing ( $\alpha_{1}, \beta_{1}$ ) by ( $\alpha_{2}, \beta_{2}$ ) (recall that $\beta_{1}=\alpha_{2}$ ), we obtain a map $v_{2} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\beta_{2}, N\right\}, S^{2}\right) \cap H^{1}(\Omega)$ satisfying $v_{2}=v_{1}$ outside a $\delta$-neighborhood of $\left[\alpha_{2}, \beta_{2}\right]$ included in $\Omega, \operatorname{deg}\left(v_{2}, \beta_{2}\right)=+1, \operatorname{deg}\left(v_{2}, N\right)=-1$ and

$$
\int_{\Omega}\left|\nabla v_{2}(x)\right|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla v_{1}(x)\right|^{2} w(x) d x+8 \liminf _{m \rightarrow+\infty} \int_{\Theta_{m}\left(\left[\alpha_{2}, \beta_{2}\right]\right)} \varepsilon_{\alpha_{2}, \beta_{2}}^{m}(x) w(x) d x+C \delta .
$$

Iterating this process, we finally obtain a map $v_{n-1} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\alpha_{n}, \beta_{n}\right\}, S^{2}\right) \cap H^{1}(\Omega)$ (recall that $\left.\beta_{n}=N\right)$ verifying $v_{n-1}=u$ outside a $\delta$-neighborhood of $\cup_{k=1}^{n-1}\left[\alpha_{k}, \beta_{k}\right]$ included in $\Omega$, $\operatorname{deg}\left(v_{n-1}, \alpha_{n}\right)=+1, \operatorname{deg}\left(v_{n-1}, \beta_{n}\right)=-1$ and

$$
\int_{\Omega}\left|\nabla v_{n-1}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \sum_{k=1}^{n-1} \liminf _{m \rightarrow+\infty} \int_{\Theta_{m}\left(\left[\alpha_{k}, \beta_{k}\right]\right)} \varepsilon_{\alpha_{k}, \beta_{k}}^{m}(x) w(x) d x+C \delta .
$$

As in Step 1, we consider $0<r_{0}<\delta$ such that $B_{r_{0}}\left(\alpha_{n}\right) \cap B_{r_{0}}\left(\beta_{n}\right)=\emptyset$ and $B_{r_{0}}\left(\alpha_{n}\right) \cup B_{r_{0}}\left(\beta_{n}\right) \subset \Omega$ and we construct, using Lemma A1 in [1], a map $\tilde{v} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\alpha_{n}, \beta_{n}\right\}, S^{2}\right) \cap H^{1}(\Omega)$ satisfying

$$
\tilde{v}(x)= \begin{cases}u(x) & \text { on } \Omega \backslash B_{r_{0}}\left(\alpha_{n}\right), \\ R_{+}\left(\frac{x-\alpha_{n}}{\left|x-\alpha_{n}\right|}\right) & \text { on } B_{r_{0}}\left(\alpha_{n}\right), \\ -R_{-}\left(\frac{x-\beta_{n}}{\left|x-\beta_{n}\right|}\right) & \text { on } B_{r_{0}}\left(\beta_{n}\right),\end{cases}
$$

for some rotations $R_{+}$and $R_{-}$and

$$
\int_{\Omega}|\nabla \tilde{v}(x)|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla v_{n-1}(x)\right|^{2} w(x) d x+\delta .
$$

Applying the construction in Step 2 starting from $\tilde{v}$, we obtain a new map $\tilde{v}_{n}^{m_{n}}$ (for some large $m_{n} \in \mathbb{N}$ ) defined on $\delta$-neighborhood $W^{\prime}$ of $\left[\alpha_{n}, \beta_{n}\right]$ included in $\Omega$, which coincide with $\tilde{v}$ near $\partial W^{\prime}$, which then has only point singularities of degree zero (since $\operatorname{deg}\left(\tilde{v}, \beta_{n}\right)=-1$ ) and satisfying

$$
\int_{W^{\prime}}\left|\nabla v_{n}^{m_{n}}(x)\right|^{2} w(x) d x \leq \int_{W^{\prime}}|\nabla \tilde{v}(x)|^{2} w(x) d x+8 \liminf _{m \rightarrow+\infty} \int_{\Theta_{m}\left(\left[\alpha_{n}, \beta_{n}\right]\right)} \varepsilon_{\alpha_{n}, \beta_{n}}^{m}(x) w(x) d x+C \delta .
$$

Since the degree of each singularities of $v_{n}^{m_{n}}$ is zero, we can construct a map $v_{n} \in \mathcal{C}^{1}\left(\bar{W}^{\prime}, S^{2}\right)$ (see $[2,3]$ ) verifying $v_{n}=\tilde{v}$ in a neighborhood of $\partial W^{\prime}$ and

$$
\int_{W^{\prime}}\left|\nabla v_{n}(x)\right|^{2} w(x) d x \leq \int_{W^{\prime}}\left|\nabla v_{n}^{m_{n}}(x)\right|^{2} w(x) d x+\delta
$$

Then we define $u_{\delta}: \bar{\Omega} \rightarrow S^{2}$ by

$$
u_{\delta}(x)= \begin{cases}v_{n-1}(x) & \text { if } x \in \bar{\Omega} \backslash W^{\prime} \\ v_{n}(x) & \text { if } x \in \bar{W}^{\prime}\end{cases}
$$

Since $v_{n-1}=\tilde{v}$ and $\tilde{v}=v_{n-1}$ near $\partial W^{\prime}$, we deduce that $u_{\delta} \in \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$. Moreover it follows by construction that $u_{\delta}=u$ outside a $\delta$-neighborhood of $\cup_{k=1}^{n}\left[\alpha_{k}, \beta_{k}\right]$ included in $\Omega$ and

$$
\int_{\Omega}\left|\nabla u_{\delta}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \bar{\ell}(\mathcal{F})+C \delta,
$$

which ends the proof since $\delta$ is arbitrary small.

Lemma 2.3. Let $\left(P_{i}, N_{i}\right)_{i=1}^{K}$ be $2 K$ distinct points in $\Omega$ and consider $u \in H^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega} \backslash$ $\left.\cup_{i=1}^{K}\left\{P_{i}, N_{i}\right\}\right)$ such that $\operatorname{deg}\left(u, P_{i}\right)=+1$ and $\operatorname{deg}\left(u, N_{i}\right)=-1$ for $i=1, \ldots, K$. Then There exists a sequence of maps $\left(u_{n}\right)_{n \in \mathbb{N}} \subset \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ satisfying $u_{n \mid \partial \Omega}=u_{\mid \partial \Omega}$,

$$
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u)+2^{-n},
$$

and

$$
\text { meas }\left(\left\{x \in \Omega, u_{n}(x) \neq u(x)\right\}\right) \leq 2^{-n} .
$$

Proof. Without loss of generality we may assume that $\sum_{i} d_{w}\left(P_{i}, N_{i}\right)$ is equal to the length of a minimal connection relative to $d_{w}$ between the points $\left(P_{i}\right)$ and $\left(N_{i}\right)$. As in [1], we are going to "remove" each dipole ( $P_{i}, N_{i}$ ). More precisely, for each $n \in \mathbb{N}$, we construct successively $K$ maps $\left(u_{n}^{i}\right)_{i=1}^{K}$ satisfying
(a) $u_{n}^{i} \in H^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}\left(\bar{\Omega} \backslash \bigcup_{i+1 \leq j \leq K}\left\{P_{j}, N_{j}\right\}\right)$ for $i=1, \ldots, K$,
(b) $u_{n}^{1}=u$ on $\bar{\Omega} \backslash W_{n}^{1}$ and $u_{n}^{i}=u_{n}^{i-1}$ on $\bar{\Omega} \backslash W_{n}^{i}$ for $i=2, \ldots, K$ where $W_{n}^{i}$ is is strictly included in $\Omega \backslash \bigcup_{i+1 \leq j \leq K}\left\{P_{j}, N_{j}\right\}$ and $\left|W_{n}^{i}\right| \leq 2^{-n} / K$,
(c) $\int_{\Omega}\left|\nabla u_{n}^{1}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi d_{w}\left(P_{1}, N_{1}\right)+\frac{2^{-n}}{K}$ and

$$
\int_{\Omega}\left|\nabla u_{n}^{i}(x)\right|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla u_{n}^{i-1}(x)\right|^{2} w(x) d x+8 \pi d_{w}\left(P_{i}, N_{i}\right)+\frac{2^{-n}}{K} \text { for } i=2, \ldots, K
$$

We easily check that the sequence $\left(u_{n}^{K}\right)_{n \in \mathbb{N}}$ then satisfies the requirement since we have $L_{w}(u)=\sum_{i} d_{w}\left(P_{i}, N_{i}\right)$. We start with the construction of $u_{n}^{1}$.
Construction of $u_{n}^{1}$. By Lemma 2.1, we can find $\mathcal{F}_{1}=\left(\left[\alpha_{1}, \beta_{1}\right], \ldots,\left[\alpha_{l}, \beta_{l}\right]\right) \in \mathcal{Q}\left(P_{1}, N_{1}\right)$ satisfying

$$
\begin{equation*}
\left(\cup_{i=2}^{K}\left\{P_{i}, N_{i}\right\} \cup\left\{N_{1}\right\}\right) \cap\left(\cup _ { k = 2 } ^ { l } [ \alpha _ { k } , \beta _ { k } ] \cup \left[\alpha_{1}, \beta_{1}[)=\emptyset,\right.\right. \tag{2.13}
\end{equation*}
$$

and

$$
\bar{\ell}_{w}\left(\mathcal{F}_{1}\right) \leq d_{w}\left(P_{1}, N_{1}\right)+\frac{2^{-(n+1)}}{8 K \pi} .
$$

From (2.13), we infer that we can find $\delta>0$ small enough such that

$$
W_{\delta}^{1}=\left\{x \in \mathbb{R}^{3}, \operatorname{dist}\left(x, \cup_{k=1}^{l}\left[\alpha_{k}, \beta_{k}\right]\right) \leq \delta\right\} \subset \Omega \backslash \cup_{i=2}^{K}\left\{P_{i}, N_{i}\right\} \quad \text { and } \quad\left|W_{\delta}^{1}\right| \leq \frac{2^{-n}}{K}
$$

By the method described in the proof of Lemma 2.2, we construct a map $u_{n}^{1} \in H^{1}\left(\Omega, S^{2}\right) \cap$ $\mathcal{C}^{1}\left(\bar{\Omega} \backslash \cup_{i=2}^{K}\left\{P_{i}, N_{i}\right\}\right)$ verifying $u_{n}^{1}=u$ outside $W_{\delta}^{1}$ and

$$
\begin{aligned}
\int_{\Omega}\left|\nabla u_{n}^{1}(x)\right|^{2} w(x) d x a m p & ; \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \bar{\ell}_{w}\left(\mathcal{F}_{1}\right)+\frac{2^{-(n+1)}}{K} \\
a m p & ; \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi d_{w}\left(P_{1}, N_{1}\right)+\frac{2^{-n}}{K} .
\end{aligned}
$$

Construction of $u_{n}^{i}, i=2, \ldots, K$. We iterate the previous process i.e., we proceed as for the construction of $u_{n}^{1}$ but starting from $u_{n}^{i-1}$ instead of $u$.

## 3 Proof of Theorem 1.2

### 3.1 Lower Bound of the Energy

In this section, we denote by $F_{w}$ the functional defined for maps $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ by

$$
F_{w}(u)=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u) .
$$

Proposition 3.1. The functional $F_{w}$ is sequentially lower semi-continuous on $H_{g}^{1}\left(\Omega, S^{2}\right)$ for the weak $H^{1}$-topology.

Proof. We follow the method in [4]. Since the supremum of a family of sequentially lower semi-continuous functionals is sequentially lower semi-continuous, it suffices to show that for any function $\zeta: \bar{\Omega} \rightarrow \mathbb{R}$ which is 1 -Lipschitz with respect to $d_{w}$, the functional

$$
u \in H_{g}^{1} \mapsto \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+2 \int_{\Omega} D(u) \cdot \nabla \zeta d x
$$

is sequentially lower semi-continuous for the weak $H^{1}$-topology (the term $\int_{\partial \Omega}(D(u) \cdot \nu) \zeta$ only depends on $g$ and $\zeta)$. Consider $\left(u_{n}\right)_{n \in \mathbb{N}} \subset H_{g}^{1}\left(\Omega, S^{2}\right)$ and $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ such that $u_{n} \rightharpoonup u$ weakly in $H^{1}$. Setting $v_{n}=u_{n}-u$, we have

$$
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+\int_{\Omega}\left|\nabla v_{n}(x)\right|^{2} w(x) d x+o(1)
$$

and writing

$$
2 \int_{\Omega} D\left(u_{n}\right) \cdot \nabla \zeta d x=A_{n}+B_{n}+C_{n}
$$

with

$$
\begin{aligned}
A_{n}= & 2 \int_{\Omega} u_{n} \cdot\left(\frac{\partial u}{\partial x_{2}} \wedge \frac{\partial u}{\partial x_{3}} \frac{\partial \zeta}{\partial x_{1}}+\frac{\partial u}{\partial x_{3}} \wedge \frac{\partial u}{\partial x_{1}} \frac{\partial \zeta}{\partial x_{3}}+\frac{\partial u}{\partial x_{1}} \wedge \frac{\partial u}{\partial x_{2}} \frac{\partial \zeta}{\partial x_{3}}\right) \\
B_{n}= & 2 \int_{\Omega} u_{n} \cdot\left(\frac{\partial v_{n}}{\partial x_{2}} \wedge \frac{\partial u}{\partial x_{3}}+\frac{\partial u}{\partial x_{2}} \wedge \frac{\partial v_{n}}{\partial x_{3}}\right) \frac{\partial \zeta}{\partial x_{1}}+2 \int_{\Omega} u_{n} \cdot\left(\frac{\partial v_{n}}{\partial x_{3}} \wedge \frac{\partial u}{\partial x_{1}}+\frac{\partial u}{\partial x_{3}} \wedge \frac{\partial v_{n}}{\partial x_{1}}\right) \frac{\partial \zeta}{\partial x_{2}} \\
& +2 \int_{\Omega} u_{n} \cdot\left(\frac{\partial v_{n}}{\partial x_{1}} \wedge \frac{\partial u}{\partial x_{2}}+\frac{\partial u}{\partial x_{1}} \wedge \frac{\partial v_{n}}{\partial x_{2}}\right) \frac{\partial \zeta}{\partial x_{3}} \\
C_{n}= & 2 \int_{\Omega} u_{n} \cdot\left(\frac{\partial v_{n}}{\partial x_{2}} \wedge \frac{\partial v_{n}}{\partial x_{3}} \frac{\partial \zeta}{\partial x_{1}}+\frac{\partial v_{n}}{\partial x_{3}} \wedge \frac{\partial v_{n}}{\partial x_{1}} \frac{\partial \zeta}{\partial x_{3}}+\frac{\partial v_{n}}{\partial x_{1}} \wedge \frac{\partial v_{n}}{\partial x_{2}} \frac{\partial \zeta}{\partial x_{3}}\right)
\end{aligned}
$$

We easily obtain that $A_{n} \rightarrow 2 \int_{\Omega} D(u) \cdot \nabla \zeta$ as $n \rightarrow+\infty$ since $u_{n} \rightharpoonup u$ weak» in $L^{\infty}$ and that $B_{n} \rightarrow 0$ since $v_{n} \rightharpoonup 0$ weakly in $L^{2}$ and $u_{n} \rightarrow u$ strongly in $L^{2}$. Now we set

$$
V_{n}=\left(u_{n} \cdot \frac{\partial v_{n}}{\partial x_{2}} \wedge \frac{\partial v_{n}}{\partial x_{3}}, u_{n} \cdot \frac{\partial v_{n}}{\partial x_{3}} \wedge \frac{\partial v_{n}}{\partial x_{1}}, u_{n} \cdot \frac{\partial v_{n}}{\partial x_{1}} \wedge \frac{\partial v_{n}}{\partial x_{2}}\right)
$$

We have

$$
\left|C_{n}\right|=2\left|\int_{\Omega} V_{n} \cdot \nabla \zeta\right| \leq 2 \int_{\Omega}\left|V_{n}\right||\nabla \zeta| .
$$

By Lemma 1 in [4], we know that $2\left|V_{n}\right| \leq\left|\nabla v_{n}\right|^{2}$ and by Proposition 2.3 in [14], any $\zeta: \bar{\Omega} \rightarrow \mathbb{R}$ which 1-Lipschitz with respect to $d_{w}$ satisfies $|\nabla \zeta| \leq w$ a.e. on $\Omega$. Then we obtain

$$
\left|C_{n}\right| \leq \int_{\Omega}\left|\nabla v_{n}(x)\right|^{2} w(x) d x
$$

and we conclude that

$$
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x+2 \int_{\Omega} D\left(u_{n}\right) \cdot \nabla \zeta d x \geq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+2 \int_{\Omega} D(u) \cdot \nabla \zeta d x+o(1)
$$

which clearly implies the result.

Proof of " $\geq$ " in Theorem 1.2. Let $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ and consider an arbitrary sequence $\left(u_{n}\right)_{n \in \mathbb{N}} \subset H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega})$ such that $u_{n} \rightharpoonup u$ weakly in $H^{1}$. Since $u_{n}$ is smooth in $\Omega$, we have $T\left(u_{n}\right) \equiv 0$ and then $L_{w}\left(u_{n}\right)=0$. We conclude by Proposition 3.1 that

$$
\liminf _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x=\liminf _{n \rightarrow+\infty} F_{w}\left(u_{n}\right) \geq F_{w}(u)=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u)
$$

Since the sequence $\left(u_{n}\right)_{n \in \mathbb{N}}$ is arbitrary, we get the announced result.

### 3.2 Upper Bound of the Energy

Proposition 3.2. Let $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$. Then there exists a sequence of maps $\left(u_{n}\right)_{n \in \mathbb{N}} \subset$ $H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega})$ such that $u_{n} \rightharpoonup u$ weakly in $H^{1}$ and

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u) .
$$

End of the proof of Theorem 1.2. Let $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$ and let $\left(u_{n}\right)_{n \in \mathbb{N}}$ be the sequence of maps given by Proposition 3.2. By definition of $E_{w}(u)$ and Proposition 3.2, we have

$$
E_{w}(u) \leq \liminf _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u)
$$

which ends the proof of Theorem 1.2.

To prove Proposition 3.2, we need the following Lemma. We postpone its proof at the end of this section.

Lemma 3.1. For any $u, v \in H_{g}^{1}\left(\Omega, S^{2}\right)$, we have

$$
\begin{equation*}
\left|L_{w}(u)-L_{w}(v)\right| \leq C \Lambda\left(\|\nabla u\|_{L^{2}(\Omega)}+\|\nabla v\|_{L^{2}(\Omega)}\right)\|\nabla u-\nabla v\|_{L^{2}(\Omega)} \tag{3.1}
\end{equation*}
$$

for a constant $C$ independent of $w$.
Proof of Proposition 3.2. Let $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$. By the result in [1, 3], we can find a sequence of $\operatorname{maps}\left(v_{n}\right)_{n \in \mathbb{N}} \subset H_{g}^{1}\left(\Omega, S^{2}\right)$ such that $v_{n} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash \cup_{i=1}^{K_{n}}\left\{P_{i}, N_{i}\right\}\right)$ for some $2 K_{n}$ distinct points $\left(P_{i}, N_{i}\right)$ in $\Omega, \operatorname{deg}\left(v_{n}, P_{i}\right)=+1$ and $\operatorname{deg}\left(v_{n}, N_{i}\right)=-1$ for $i=1, \ldots, K_{n}$ and such that

$$
\begin{equation*}
\left\|\nabla\left(v_{n}-u\right)\right\|_{L^{2}(\Omega)} \leq 2^{-n} \tag{3.2}
\end{equation*}
$$

From this inequality we infer that

$$
\begin{equation*}
\operatorname{meas}\left(\left\{x \in \Omega,\left|v_{n}(x)-u(x)\right|<2^{-n / 2}\right\}\right) \leq C 2^{-n} \tag{3.3}
\end{equation*}
$$

Applying Lemma 2.3 to $v_{n}$, we find a map $u_{n} \in \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ satisfying $u_{n \mid \partial \Omega}=g$,

$$
\begin{equation*}
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla v_{n}(x)\right|^{2} w(x) d x+8 \pi L_{w}\left(v_{n}\right)+2^{-n} \tag{3.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\text { meas }\left(\left\{x \in \Omega, u_{n}(x) \neq v_{n}(x)\right\}\right) \leq 2^{-n} . \tag{3.5}
\end{equation*}
$$

From (3.2) and Lemma 3.1 we deduce that $L_{w}\left(v_{n}\right) \rightarrow L_{w}(u)$ as $n \rightarrow+\infty$ and then it follows that $\left(u_{n}\right)_{n \in \mathbb{N}}$ is bounded in $H^{1}$. Moreover we obtain from (3.3) and (3.5) that $u_{n} \rightarrow u$ a.e. in $\Omega$ and we conclude that $u_{n} \rightharpoonup u$ weakly in $H^{1}$. Letting $n \rightarrow+\infty$ in (3.4) leads to

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi L_{w}(u)
$$

which completes the proof.
Proof of Lemma 3.1. To prove Lemma 3.1, we follow the method in [4]. For $u, v \in H_{g}^{1}\left(\Omega, S^{2}\right)$, we set

$$
L_{w}(u, v)=\operatorname{Sup}\left\{\int_{\Omega}(D(u)-D(v)) \cdot \nabla \zeta, \zeta: \bar{\Omega} \rightarrow \mathbb{R} 1 \text {-Lipschitz with respect to } d_{w}\right\} .
$$

Since $D(u) \cdot \nu=D(v) \cdot \nu$ on $\partial \Omega$ (it only depends on $g$ ), we have

$$
\int_{\Omega} D(u) \cdot \nabla \zeta-\int_{\partial \Omega}(D(u) \cdot \nu) \zeta=\int_{\Omega} D(v) \cdot \nabla \zeta-\int_{\partial \Omega}(D(v) \cdot \nu) \zeta+\int_{\Omega}(D(u)-D(v)) \cdot \nabla \zeta,
$$

and we easily derive that

$$
\left|L_{w}(u)-L_{w}(v)\right| \leq L_{w}(u, v) .
$$

Similar computations to those in [4], proof of Theorem 1, lead to

$$
\left|\int_{\Omega}(D(u)-D(v)) \cdot \nabla \zeta\right| \leq C\left(\|\nabla u\|_{L^{2}(\Omega)}+\|\nabla v\|_{L^{2}(\Omega)}\right)\|\nabla u-\nabla v\|_{L^{2}(\Omega)}\|\nabla \zeta\|_{L^{\infty}(\Omega)} .
$$

By Proposition 2.3 in [14], any real function $\zeta$ which is 1-Lipschitz with respect to $d_{w}$ satisfies $|\nabla \zeta| \leq w$ a.e. on $\Omega$. We deduce that (3.1) holds since $w \leq \Lambda$ a.e. on $\Omega$.

## 4 Stability and Approximation Properties

### 4.1 A Stability Property

Before stating the result, we need to recall some previous ones obtained in [14]. For any real measurable function $w$ satisfying assumption (1.1), we may associate to distance $d_{w}$ the length functional $\mathbb{L}_{d_{w}}$ defined by

$$
\mathbb{L}_{d_{w}}(\gamma)=\operatorname{Sup}\left\{\sum_{k=0}^{m-1} d_{w}\left(\gamma\left(t_{k}\right), \gamma\left(t_{k+1}\right)\right), 0=t_{0}<t_{1}<\ldots<t_{m}=1, m \in \mathbb{N}^{\star}\right\},
$$

where $\gamma:[0,1] \rightarrow \bar{\Omega}$ is any continuous curve. In [14], we have proved that for any $x, y \in \bar{\Omega}$,

$$
\begin{equation*}
d_{w}(x, y)=\operatorname{Inf}\left\{\mathbb{L}_{d_{w}}(\gamma), \gamma \in \operatorname{Lip}([0,1], \bar{\Omega}), \gamma(0)=x \text { and } \gamma(1)=y\right\} \tag{4.1}
\end{equation*}
$$

where $\operatorname{Lip}([0,1], \bar{\Omega})$ denotes the class of all Lipschitz maps from $[0,1]$ into $\bar{\Omega}$. We have also shown that the infimum in (4.1) is in fact achieved.

The following stability result relies on the $\Gamma$-convergence of the length functionals (we refer to [11] for the notion of $\Gamma$-convergence). In the sequel, we endow $\operatorname{Lip}([0,1], \bar{\Omega})$ with the topology of the uniform convergence on $[0,1]$.

Theorem 4.1. Let $\left(w_{n}\right)_{n \in \mathbb{N}}$ be a sequence of measurable real functions such that

$$
\begin{equation*}
0<c_{0} \leq w_{n} \leq C_{0} \quad \text { a.e. in } \Omega \tag{4.2}
\end{equation*}
$$

for some constants $c_{0}$ and $C_{0}$ independent of $n \in \mathbb{N}$. Then the following properties are equivalent:
(i) the functionals $\mathbb{L}_{d_{w_{n}}} \Gamma$-converge to $\mathbb{L}_{d_{w}}$ in $\operatorname{Lip}([0,1], \bar{\Omega})$ and

$$
\begin{equation*}
\int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x \underset{n \rightarrow+\infty}{\rightarrow} \int_{\Omega}|\nabla \varphi(x)|^{2} w(x) d x \quad \text { for any } \varphi \in H^{1}(\Omega, \mathbb{R}) \tag{4.3}
\end{equation*}
$$

(ii) for every smooth boundary data $g: \partial \Omega \rightarrow S^{2}$ such that $\operatorname{deg}(g)=0$,

$$
E_{w_{n}}(u) \underset{n \rightarrow+\infty}{\rightarrow} E_{w}(u) \quad \text { for any } u \in H_{g}^{1}\left(\Omega, S^{2}\right)
$$

Proof. (i) $\Rightarrow$ (ii). We fix a smooth boundary data $g: \Omega \rightarrow S^{2}$ such that $\operatorname{deg}(g)=0$. Clearly (4.3) implies that

$$
\int_{\Omega}|\nabla u(x)|^{2} w_{n}(x) d x \underset{n \rightarrow+\infty}{\rightarrow} \int_{\Omega}|\nabla u(x)|^{2} w(x) d x \quad \text { for any } u \in H_{g}^{1}\left(\Omega, S^{2}\right),
$$

and by Theorem 1.2, it remains to prove that

$$
\begin{equation*}
L_{w_{n}}(u) \underset{n \rightarrow+\infty}{\rightarrow} L_{w}(u) \text { for any } u \in H_{g}^{1}\left(\Omega, S^{2}\right) \tag{4.4}
\end{equation*}
$$

Consider $u \in H_{g}^{1}\left(\Omega, S^{2}\right)$. By the result in [1, 3], there exits a sequence of maps $\left(v_{k}\right)_{k \in \mathbb{N}} \subset$ $H_{g}^{1}\left(\Omega, S^{2}\right)$ such that $v_{k} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash \cup_{j=1}^{M_{k}}\left\{P_{j}, N_{j}\right\}, S^{2}\right)$ for some $2 M_{k}$ points ( $P_{j}, N_{j}$ ) in $\Omega$, $\operatorname{deg}\left(v_{k}, P_{j}\right)=+1$ and $\operatorname{deg}\left(v_{k}, N_{j}\right)=-1$ for $j=1, \ldots, M_{k}$, and $v_{k} \rightarrow u$ strongly in $H^{1}$. We have

$$
L_{w_{n}}\left(v_{k}\right)=\operatorname{Min}_{\sigma \in \mathcal{S}_{M_{k}}} \sum_{j=1}^{M_{k}} d_{w_{n}}\left(P_{j}, N_{\sigma(j)}\right) \quad \text { and } \quad L_{w}\left(v_{k}\right)=\operatorname{Min}_{\sigma \in \mathcal{S}_{M_{k}}} \sum_{j=1}^{M_{k}} d_{w}\left(P_{j}, N_{\sigma(j)}\right)
$$

Since the functionals $\mathbb{L}_{d_{w_{n}}} \Gamma$-converge to $\mathbb{L}_{d_{w}}$ in $\operatorname{Lip}([0,1], \bar{\Omega})$, we deduce from Theorem 4.1 in [14] that for every $k \in \mathbb{N}, L_{w_{n}}\left(v_{k}\right) \rightarrow L_{w}\left(v_{k}\right)$ as $n \rightarrow+\infty$. Now we fix a small $\delta>0$. Since $v_{k} \rightarrow u$ strongly in $H^{1}$, we derive from Lemma 3.1 and (4.2) that exists $k_{0} \in \mathbb{N}$ which only depends on $u, \delta$ and $C_{0}$ such that

$$
L_{w_{n}}\left(v_{k}\right)-\delta \leq L_{w_{n}}(u) \leq L_{w_{n}}\left(v_{k}\right)+\delta \quad \text { for any } n \in \mathbb{N} \text { and } k \geq k_{0} .
$$

Letting $n \rightarrow+\infty$ in this inequality, we get that

$$
L_{w}\left(v_{k}\right)-\delta \leq \liminf _{n \rightarrow+\infty} L_{w_{n}}(u) \leq \limsup _{n \rightarrow+\infty} L_{w_{n}}(u) \leq L_{w}\left(v_{k}\right)+\delta \quad \text { for } k \geq k_{0} .
$$

Passing to the limit in $k$ and using Lemma 3.1, we obtain

$$
L_{w}(u)-\delta \leq \liminf _{n \rightarrow+\infty} L_{w_{n}}(u) \leq \limsup _{n \rightarrow+\infty} L_{w_{n}}(u) \leq L_{w}(u)+\delta,
$$

which leads to the result since $\delta$ is arbitrary small.
(ii) $\Rightarrow$ (i). First we prove (4.3) for $\varphi \in \mathcal{C}^{\infty}(\bar{\Omega}, \mathbb{R})$. Let $\varphi \in \mathcal{C}^{\infty}(\bar{\Omega}, \mathbb{R})$ and consider the smooth map $g: \partial \Omega \rightarrow S^{2}$ defined by $g(x)=(\cos (\varphi(x)), \sin (\varphi(x)), 0)$. We easily check that $\operatorname{deg}(g)=0$. Now consider the map $u$ defined for $x \in \bar{\Omega}$ by

$$
u(x)=(\cos (\varphi(x)), \sin (\varphi(x)), 0) .
$$

We have $u \in H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{\infty}(\bar{\Omega})$ and then $L_{w_{n}}(u)=L_{w}(u)=0$ for any $n \in \mathbb{N}$. Since $|\nabla u|^{2}=|\nabla \varphi|^{2}$, we derive from assumption (ii) and Theorem 1.2 that

$$
\int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x \underset{n \rightarrow+\infty}{\rightarrow} \int_{\Omega}|\nabla \varphi(x)|^{2} w(x) d x
$$

Let us now prove (4.3) for any $\varphi \in H^{1}(\Omega, \mathbb{R})$. Let $\varphi \in H^{1}(\Omega, \mathbb{R})$ and consider a sequence $\left(\varphi_{k}\right)_{k \in \mathbb{N}} \subset \mathcal{C}^{\infty}(\bar{\Omega}, \mathbb{R})$ such that $\varphi_{k} \rightarrow \varphi$ strongly in $H^{1}$. We fix a small $\delta>0$. From assumption (4.2), we infer that exists $k_{0} \in \mathbb{N}$ which only depends on $\varphi, \delta$ and $C_{0}$ such that for any $n \in \mathbb{N}$ and $k \geq k_{0}$,

$$
\int_{\Omega}\left|\nabla \varphi_{k}(x)\right|^{2} w_{n}(x) d x-\delta \leq \int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x \leq \int_{\Omega}\left|\nabla \varphi_{k}(x)\right|^{2} w_{n}(x) d x+\delta .
$$

Since $\varphi_{k}$ is smooth, letting $n \rightarrow+\infty$ we obtain for $k \geq k_{0}$,

$$
\begin{aligned}
\int_{\Omega}\left|\nabla \varphi_{k}(x)\right|^{2} w(x) d x-\delta & \leq \liminf _{n \rightarrow+\infty} \int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x \\
& \leq \limsup _{n \rightarrow+\infty} \int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x \leq \int_{\Omega}\left|\nabla \varphi_{k}(x)\right|^{2} w(x) d x+\delta .
\end{aligned}
$$

Passing to the limit in $k$ and then $\delta \rightarrow 0$, we conclude

$$
\lim _{n \rightarrow+\infty} \int_{\Omega}|\nabla \varphi(x)|^{2} w_{n}(x) d x=\int_{\Omega}|\nabla \varphi(x)|^{2} w(x) d x .
$$

It remains to prove that the functionals $\mathbb{L}_{d_{w_{n}}} \Gamma$-converge to $\mathbb{L}_{d_{w}}$ in $\operatorname{Lip}([0,1], \bar{\Omega})$. Let $P$ and $N$ be two distinct points in $\Omega$. We take $g \equiv(0,0,1)$ and consider $u \in H_{g}^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}(\bar{\Omega} \backslash\{P, N\})$ (such a map is constructed for instance in $[6,9]$ ). By Theorem 1.2, we have

$$
E_{w_{n}}(u)=\int_{\Omega}|\nabla u(x)|^{2} w_{n}(x) d x+8 \pi d_{w_{n}}(P, N)
$$

and

$$
E_{w}(u)=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi d_{w}(P, N)
$$

From (4.3) we get that $\int_{\Omega}|\nabla u(x)|^{2} w_{n}(x) d x \rightarrow \int_{\Omega}|\nabla u(x)|^{2} w(x) d x$ and from assumption (ii) we deduce that

$$
d_{w_{n}}(P, N) \rightarrow d_{w}(P, N) \quad \text { as } n \rightarrow+\infty
$$

Since the points $P$ and $N$ are arbitrary in $\Omega$, we derive that $d_{w_{n}}$ converges to $d_{w}$ pointwise on $\Omega \times \Omega$ and the conclusion follows by the results in [14] Section 4 .

In the next proposition, we give some sufficient condition on a sequence $\left(w_{n}\right)_{n \in \mathbb{N}}$ converging pointwise a.e. to $w$ for property (ii) in Theorem 4.1 to hold.

Proposition 4.1. Let $\left(w_{n}\right)_{n \in \mathbb{N}}$ be a sequence of measurable real functions satisfying (4.2) and assume that one of the following conditions holds:
(a) $w_{n} \geq w$ and $w_{n} \rightarrow w$ a.e. in $\Omega$,
(b) $w_{n} \rightarrow w$ in $L^{\infty}(\Omega)$.

Then property (ii) in Theorem 4.1 holds.
Proof. By Proposition 4.1 and Theorem 4.1 in [14], (a) or (b) implies that the functionals $\mathbb{L}_{d_{w_{n}}} \Gamma$-converge to $\mathbb{L}_{d_{w}}$ in $\operatorname{Lip}([0,1], \bar{\Omega})$. We also check that (a) or (b) implies (4.3) by dominated convergence. Then the conclusion follows from Theorem 4.1.

Remark 4.1. The conclusion of Proposition 4.1 may fails if one only assumes that $w_{n} \rightarrow w$ a.e. in $\Omega$ (see Remark 4.1 in [14]).

### 4.2 Approximation Property

In this section, we show that the functional $E_{w}$ can be obtain as pointwise limit of a sequence $\left(E_{w_{n}}\right)_{n \in \mathbb{N}}$ in which the weight function $w_{n}$ is smooth.

Proposition 4.2. Let $\left(\rho_{n}\right)_{n \in \mathbb{N}}$ be a sequence of smooth mollifiers. Extending $w$ by a sufficiently large constant and setting $w_{n}=\rho_{n} * w$, we have

$$
E_{w_{n}}(u) \underset{n \rightarrow+\infty}{\rightarrow} E_{w}(u) \quad \text { for any } u \in H_{g}^{1}\left(\Omega, S^{2}\right) .
$$

Proof. By construction, (4.3) clearly holds. Then property (i) in Theorem 4.1 follows from Theorem 4.1 in [14] and Theorem 4.2 in [14] which leads to the result by Theorem 4.1.

## 5 The Relaxed Energy without Prescribed Boundary Data

In this section, we consider the relaxed type functional

$$
\tilde{E}_{w}(u)=\operatorname{Inf}\left\{\liminf _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x, u_{n} \in \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right), u_{n} \rightharpoonup u \text { weakly in } H^{1}\right\}
$$

defined for $u \in H^{1}\left(\Omega, S^{2}\right)$. We recall that F . Bethuel has also proved (see [1]) that $\mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ is sequentially dense in $H^{1}\left(\Omega, S^{2}\right)$ for the weak $H^{1}$ topology and then $\tilde{E}_{w}$ is well defined.

As in [4], there is also a notion of length of a minimal connection relative to $d_{w}$ defined for any $u \in H^{1}\left(\Omega, S^{2}\right)$ :

$$
\tilde{L}_{w}(u)=\frac{1}{4 \pi} \operatorname{Sup}\left\{\langle T(u), \zeta\rangle, \zeta: \bar{\Omega} \rightarrow \mathbb{R} 1 \text {-Lipschitz with respect to } d_{w} \text { and } \zeta=0 \text { on } \partial \Omega\right\}
$$

Since no assumptions are made on $u_{\mid \partial \Omega}$, it may happen that $\operatorname{deg}\left(u_{\mid \partial \Omega}\right) \neq 0$ or that $\operatorname{deg}\left(u_{\mid \partial \Omega}\right)$ is not well defined. But clearly $\tilde{L}_{w}(u)$ always makes sense. When $u$ is smooth except at a finite number of point in $\Omega, \tilde{L}_{w}(u)$ is equal to the length of a minimal connection relative to $d_{w}$ between the singularities of $u$ and some virtual singularities on the boundary (see [9]). More precisely, one adds some virtual singularities on the boundary in such a way that the new configuration has the same number of positive and negative points and one consider the length of a minimal connection relative to $d_{w}$ for this configuration. Then $\tilde{L}_{w}(u)$ corresponds to the infimum of these quantities when one varies the position and the number of the boundary points. There is the variant of Theorem 1.2 for $\tilde{E}_{w}$.

Theorem 5.1. For any $u \in H^{1}\left(\Omega, S^{2}\right)$, we have

$$
\tilde{E}_{w}(u)=\int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \tilde{L}_{w}(u)
$$

### 5.1 Proof of Theorem 5.1

The inequality " $\geq$ " in Theorem 5.1 can be proved using a method similar to the one used in Section 3.1 and we omit it. We obtain " $\leq$ " as in Section 3.2 using Proposition 5.1 and Lemma 5.1 below instead of Proposition 3.2 and Lemma 3.1. The proof of Lemma 5.1 is almost identical to the proof of Lemma 3.1 and we also omit it (note that all the boundary integrals vanish since $\zeta=0$ on $\partial \Omega$ ).

Proposition 5.1. Let $u \in H^{1}\left(\Omega, S^{2}\right)$. Then there exists a sequence of maps $\left(u_{n}\right)_{n \in \mathbb{N}} \subset$ $\mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ such that

$$
u_{n} \rightharpoonup u \quad \text { weakly in } H^{1}
$$

and

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \tilde{L}_{w}(u)
$$

Lemma 5.1. For any $u, v \in H^{1}\left(\Omega, S^{2}\right)$, we have

$$
\begin{equation*}
\left|\tilde{L}_{w}(u)-\tilde{L}_{w}(v)\right| \leq C \Lambda\left(\|\nabla u\|_{L^{2}(\Omega)}+\|\nabla v\|_{L^{2}(\Omega)}\right)\|\nabla u-\nabla v\|_{L^{2}(\Omega)} \tag{5.1}
\end{equation*}
$$

for a constant $C$ independent of $w$.
Proof of Proposition 5.1. Let $u \in H^{1}\left(\Omega, S^{2}\right)$. By the result in [1, 3], we can find a sequence $\left(v_{n}\right)_{n \in \mathbb{N}} \subset H^{1}\left(\Omega, S^{2}\right)$ such that $v_{n} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash\left\{\left(a_{i}\right)_{i=1}^{N_{n}}\right\}\right)$ for some $N_{n}$ distinct points $a_{1}, \ldots, a_{N_{n}}$ in $\Omega$ and

$$
\begin{equation*}
\left\|u-v_{n}\right\|_{H^{1}(\Omega)} \leq 2^{-n} \tag{5.2}
\end{equation*}
$$

Since we are working with an approximating sequence, we may assume that $\left|\operatorname{deg}\left(v_{n}, a_{i}\right)\right|=1$ for $i=1, \ldots, N_{n}$ (see [1]). Since $v_{n}$ is smooth except at a finite number of point in $\Omega$, the length of a minimal connection $\tilde{L}_{w}\left(v_{n}\right)$ is computed as follows (see [9], part II). We pair each singularity $a_{i}$ either to another singularity in $\Omega$ of opposite degree or to a virtual singularity on the boundary with opposite degree. In other words, we allow connections to the boundary of $\Omega$. Pairing all the singularities in this way, we take a configuration that minimizes the sum of the distances between the paired singularities, computing the distances with $d_{w}$. We relabel all the singularities (the $a_{i}$ 's and the virtual singularities on the boundary), according to their multiplicity for those on the boundary, as a list of positive and negative points say $\left(P_{1}, \ldots, P_{K_{n}}\right)$ and $\left(N_{1}, \ldots, N_{K_{n}}\right)$ such that

$$
\tilde{L}_{w}\left(v_{n}\right)=\sum_{j=1}^{K_{n}} d_{w}\left(P_{j}, N_{j}\right)
$$

Using Lemma 2 bis in [1], we can find $\tilde{v}_{n} \in H^{1}\left(\Omega, S^{2}\right) \cap \mathcal{C}^{1}\left(\bar{\Omega} \backslash \cup_{j=1}^{K_{n}}\left\{\tilde{P}_{j}, \tilde{N}_{j}\right\}\right)$ for some $2 K_{n}$ distinct points $\left(\tilde{P}_{j}, \tilde{N}_{j}\right)$ in $\Omega$ such that $\tilde{v}_{n}=v_{n}$ outside a small neighborhood of $\partial \Omega$, $\operatorname{deg}\left(\tilde{v}_{n}, \tilde{P}_{j}\right)=+1$ and $\operatorname{deg}\left(\tilde{v}_{n}, \tilde{N}_{j}\right)=-1$ for $j=1, \ldots, K_{n}, \tilde{P}_{j}=P_{j}$ (respectively $\left.\tilde{N}_{j}=N_{j}\right)$ if $P_{j} \in \Omega$ (respectively if $N_{j} \in \Omega$ ) and $\left|\tilde{P}_{j}-P_{j}\right| \leq \frac{2^{-n}}{K_{n}}$ otherwise (respectively $\left|\tilde{N}_{j}-N_{j}\right| \leq \frac{2^{-n}}{K_{n}}$ ), and

$$
\begin{equation*}
\left\|\tilde{v}_{n}-v_{n}\right\|_{H^{1}(\Omega)} \leq 2^{-n} \tag{5.3}
\end{equation*}
$$

Note that, for each pair $\left(P_{j}, N_{j}\right)$, we necessarily have $\tilde{P}_{j}=P_{j}$ or $\tilde{N}_{j}=N_{j}$ and then

$$
\begin{equation*}
\left|\sum_{j=1}^{K_{n}} d_{w}\left(P_{j}, N_{j}\right)-\sum_{j=1}^{K_{n}} d_{w}\left(\tilde{P}_{j}, \tilde{N}_{j}\right)\right| \leq C 2^{-n} \tag{5.4}
\end{equation*}
$$

and from (5.2) and (5.3), we infer that

$$
\begin{equation*}
\operatorname{meas}\left(\left\{x \in \Omega,\left|u(x)-\tilde{v}_{n}(x)\right|<2^{-n / 2}\right\}\right) \leq C 2^{-n} \text {. } \tag{5.5}
\end{equation*}
$$

Applying Lemma 2.3 to $\tilde{v}_{n}$, we find a map $u_{n} \in \mathcal{C}^{1}\left(\bar{\Omega}, S^{2}\right)$ satisfying

$$
\begin{equation*}
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla \tilde{v}_{n}(x)\right|^{2} w(x) d x+8 \pi \sum_{j=1}^{K_{n}} d_{w}\left(\tilde{P}_{j}, \tilde{N}_{j}\right)+2^{-n} \tag{5.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\operatorname{meas}\left(\left\{x \in \Omega, u_{n}(x) \neq \tilde{v}_{n}(x)\right\}\right) \leq 2^{-n} \tag{5.7}
\end{equation*}
$$

From (5.4) and (5.6), we derive that

$$
\begin{equation*}
\int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}\left|\nabla v_{n}(x)\right|^{2} w(x) d x+8 \pi \tilde{L}_{w}\left(v_{n}\right)+C 2^{-n} \tag{5.8}
\end{equation*}
$$

Since $v_{n} \rightarrow u$ strongly in $H^{1}$, we deduce from Lemma 5.1 that $\tilde{L}_{w}\left(v_{n}\right) \rightarrow \tilde{L}_{w}(u)$ as $n \rightarrow+\infty$ which implies that $\left(u_{n}\right)_{n \in \mathbb{N}}$ is bounded in $H^{1}$. From (5.3) and (5.7) we obtain $u_{n} \rightarrow u$ a.e. in $\Omega$ and then we conclude that $u_{n} \rightharpoonup u$ weakly in $H^{1}$. Passing to the limit in (5.8) leads to

$$
\limsup _{n \rightarrow+\infty} \int_{\Omega}\left|\nabla u_{n}(x)\right|^{2} w(x) d x \leq \int_{\Omega}|\nabla u(x)|^{2} w(x) d x+8 \pi \tilde{L}_{w}(u)
$$

and the proof is complete.

### 5.2 Stability and Approximation Properties for $\tilde{E}_{w}$

We present in this section the variants for $\tilde{E}_{w}$ of the results in Section 4.
Theorem 5.2. Let $\left(w_{n}\right)_{n \in \mathbb{N}}$ be a sequence of measurable real functions satisfying (4.2) and assume that (i) in Theorem 4.1 holds. Then we have

$$
\begin{equation*}
\tilde{E}_{w_{n}}(u)_{n \rightarrow+\infty}^{\rightarrow} \tilde{E}_{w}(u) \quad \text { for any } u \in H^{1}\left(\Omega, S^{2}\right) . \tag{5.9}
\end{equation*}
$$

Proof. Assumption (4.3) clearly implies that

$$
\int_{\Omega}|\nabla u(x)|^{2} w_{n}(x) d x \underset{n \rightarrow+\infty}{\rightarrow} \int_{\Omega}|\nabla u(x)|^{2} w(x) d x \quad \text { for any } u \in H^{1}\left(\Omega, S^{2}\right),
$$

and by Theorem 5.1, we just have to prove that

$$
\begin{equation*}
\tilde{L}_{w_{n}}(u) \underset{n \rightarrow+\infty}{\rightarrow} \tilde{L}_{w}(u) \text { for any } u \in H^{1}\left(\Omega, S^{2}\right) . \tag{5.10}
\end{equation*}
$$

Consider $u \in H^{1}\left(\Omega, S^{2}\right)$. By the result in [1, 3], we can find a sequence $\left(v_{k}\right)_{k \in \mathbb{N}} \subset H^{1}\left(\Omega, S^{2}\right)$ such that $v_{k} \in \mathcal{C}^{1}\left(\bar{\Omega} \backslash \cup_{i=1}^{M_{k}}\left\{a_{j}\right\}, S^{2}\right)$ for some $M_{k}$ points $\left(a_{i}\right)$ in $\Omega$ and $v_{k} \rightarrow u$ strongly in $H^{1}$. We easily check that a minimal connection for $v_{k}$ relative to distance $d_{w_{n}}$ does not allow more than $\sum_{i=1}^{M_{k}}\left|\operatorname{deg}\left(v_{k}, a_{i}\right)\right|$ connections to the boundary. Therefore, extracting a subsequence $\left(n_{l}\right)_{l \in \mathbb{N}}$, we can relabel the singularities of $v_{k}$ and the virtual singularities on the boundary given by a minimal connection relative to $d_{w_{n_{l}}}$, as a list of positive points $\left(P_{1}^{l}, \ldots, P_{K_{k}}^{l}\right)$ and a list of negative points $\left(N_{1}^{l}, \ldots, N_{K_{k}}^{l}\right)$ with $K_{k}$ independent of $l$ and such that

$$
\tilde{L}_{w_{n_{l}}}\left(v_{k}\right)=\operatorname{Min}_{\sigma \in \mathcal{S}_{K_{k}}} \sum_{j=1}^{K_{k}} d_{w_{n_{l}}}\left(P_{j}^{l}, N_{\sigma(j)}^{l}\right)=\sum_{j=1}^{K_{k}} d_{w_{n_{l}}}\left(P_{j}^{l}, N_{\sigma_{l}(j)}^{l}\right)
$$

for some permutation $\sigma_{l} \in \mathcal{S}_{K_{k}}$. Extracting another subsequence if necessary, we may assume that $\sigma_{l}=\sigma_{\star}$ is independent of $l \in \mathbb{N}$ and that $P_{j}^{l} \xrightarrow[l \rightarrow+\infty]{\rightarrow} P_{j}$ and $N_{j}^{l} \underset{l \rightarrow+\infty}{ } N_{j}$ for $j=1, \ldots, K_{k}$. From the results in [14] Section 4.1, we know that assumption (i) implies that $d_{w_{n}}$ converges to $d_{w}$ uniformly on $\bar{\Omega} \times \bar{\Omega}$ and then we have

$$
\tilde{L}_{w_{n_{l}}}\left(v_{k}\right)=\sum_{j=1}^{K_{k}} d_{w_{n_{l}}}\left(P_{j}^{l}, N_{\sigma_{\star}(j)}^{l}\right) \underset{l \rightarrow+\infty}{\rightarrow} \sum_{j=1}^{K_{k}} d_{w}\left(P_{j}, N_{\sigma_{\star}(j)}\right)
$$

By definition of $\tilde{L}_{w}\left(v_{k}\right)$, we obtain that

$$
\tilde{L}_{w}\left(v_{k}\right) \leq \lim _{l \rightarrow+\infty} \tilde{L}_{w_{n_{l}}}\left(v_{k}\right)
$$

On the other hand, we can also relabel the singularities of $v_{k}$ and the virtual singularities on the boundary given by a minimal connection relative to $d_{w}$, as a list of positive points $\left(\bar{P}_{1}, \ldots, \bar{P}_{\bar{K}}\right)$ and a list of negative points $\left(\bar{N}_{1}, \ldots, \bar{N}_{\bar{K}}\right)$ such that

$$
\tilde{L}_{w}\left(v_{k}\right)=\sum_{j=1}^{\bar{K}} d_{w}\left(\bar{P}_{j}, \bar{N}_{j}\right) .
$$

As previously, we have for any $l \in \mathbb{N}$,

$$
\tilde{L}_{w_{n_{l}}}\left(v_{k}\right) \leq \sum_{j=1}^{\bar{K}} d_{w_{n_{l}}}\left(\bar{P}_{j}, \bar{N}_{j}\right)
$$

Letting $l \rightarrow+\infty$, we obtain

$$
\lim _{l \rightarrow+\infty} \tilde{L}_{w_{n_{l}}}\left(v_{k}\right) \leq \sum_{j=1}^{\bar{K}} d_{w}\left(\bar{P}_{j}, \bar{N}_{j}\right)
$$

and then we conclude that $\lim _{l \rightarrow+\infty} \tilde{L}_{w_{n_{l}}}\left(v_{k}\right)=\tilde{L}_{w}\left(v_{k}\right)$. By uniqueness of the limit, we get that the convergence holds for the full sequence i.e.,

$$
\tilde{L}_{w_{n}}\left(v_{k}\right) \underset{n \rightarrow+\infty}{\rightarrow} \tilde{L}_{w}\left(v_{k}\right) .
$$

At this stage, we can proceed as in the proof of Theorem 4.2 (i) $\Rightarrow$ (ii) using Lemma 5.1 instead of Lemma 3.1.

We obtain the following variants of Proposition 4.1 and Proposition 4.2 using Theorem 5.2 instead of Theorem 4.1.

Proposition 5.2. Let $\left(w_{n}\right)_{n \in \mathbb{N}}$ be a sequence of measurable real functions satisfying (4.2) and assume that (a) or (b) in Proposition 4.1 holds. Then (5.9) holds.

Proposition 5.3. Let $\left(\rho_{n}\right)_{n \in \mathbb{N}}$ be a sequence of smooth mollifiers. Extending $w$ by a sufficiently large constant and setting $w_{n}=\rho_{n} * w$, then (5.9) holds.

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## References

[1] F. Bethuel, A characterization of maps in $H^{1}\left(B^{3}, S^{2}\right)$ which can be approximated by smooth maps, Ann. Inst. Henri Poincaré, Analyse Non linéaire 7 (1990), 269-286.
[2] F. Bethuel, The approximation problem for Sobolev maps between two manifolds, Acta Math. 167 (1991), 153-206.
[3] F. Bethuel, X. Zheng, Density of smooth functions between two manifolds in Sobolev spaces, J. Funct. Anal. 80 (1988), 60-75.
[4] F. Bethuel, H. Brezis, J.M. Coron, Relaxed energies for harmonic maps, in Variational Problems (H. Berestycki, J.M. Coron, I. Ekeland, eds), Birkhäuser (1990), 37-52.
[5] J. Bourgain, H. Brezis, P. Mironescu, $H^{1 / 2}$-maps with values into the circle: minimal connections, lifting, and the Ginzburg-Landau equation, to appear in Publications Math. Inst. Hautes Etudes Sci.
[6] H. Brezis, Liquid crystals and energy estimates for $S^{2}$-valued maps, in [12].
[7] H. Brezis, $S^{k}$-valued maps with singularities, in Topics in calculus of variations (Montecatini Terme 1987), Lecture Notes in Math. 1365, Springer (1989).
[8] H. Brezis, J.M. Coron, Large Solutions for Harmonic Maps in Two Dimensions, Comm. Math. Phys. 92 (1983), 203-215.
[9] H. Brezis, J.M. Coron, E. Lieb, Harmonics Maps with Defects, Comm. Math. Phys. 107 (1986), 649-705.
[10] H. Brezis, P.M. Mironescu, A.C. Ponce, $W^{1,1}$-Maps with values into $S^{1}$, in Geometric Analysis of PDE Several Complex Variables, S. Chanillo, P. Cordaro, N. Hanges, J. Hounie and A. Meziani (eds.), Contemporary Mathematics series, AMS, to appear.
[11] G. Dal Maso, Introduction to $\Gamma$-convergence, Progress in Nonlinear Differential Equations and their Applications 8, Birkhäuser (1993).
[12] J. Ericksen, D. Kinderlehrer ed., Theory and Applications of liquid crystals, IMA Series 5, Springer (1987).
[13] M. Giaquinta, G. Modica, J. Souček, Cartesian Currents in the Calculus of Variations, Springer (1998).
[14] V. Millot, Energy with Weight for $S^{2}$-Valued Maps with Prescribed Singularities, to appear in Calculus of Variations and PDEs.

